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## Key Points:

- Continued warming is explained using theory from a heat balance and a carbon inventory budget
- Peak global warming is delayed after the cessation of carbon emissions
- Timing of peak warming is controlled by the weakening of ocean heat uptake

## Supporting Information:

- Supporting Information S1

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## Drivers of Continued Surface Warming After Cessation of Carbon Emissions

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**Abstract** The climate response after cessation of carbon emissions is examined here, exploiting a single equation connecting surface warming to cumulative carbon emissions. The multicentennial response to an idealized pulse of carbon is considered by diagnosing a 1,000 year integration of an Earth system model (Geophysical Fluid Dynamics Laboratory ESM2M) and an ensemble of efficient Earth system model simulations. After emissions cease, surface temperature evolves according to (i) how much of the emitted carbon remains in the atmosphere and (ii) how much of the additional radiative forcing warms the surface rather than the ocean interior. The peak in surface temperature is delayed in time after carbon emissions cease through the decline in ocean heat uptake, which in turn increases the proportion of radiative forcing warming the surface. Eventually, after many centuries, surface temperature declines as the radiative forcing decreases through the excess atmospheric CO<sub>2</sub> being taken up by the ocean and land.

**Plain Language Summary** The climate response after carbon emissions cease is examined here, exploiting a combination of theory and diagnostics of an Earth system model. The multicentennial response to an idealized pulse of carbon is considered over 1,000 years. After emissions cease, surface temperature evolves according to (i) how much of the emitted carbon remains in the atmosphere and (ii) how much of the additional radiative forcing warms the surface rather than the ocean interior. Surface temperature continues to increase after carbon emissions cease through a decline in ocean heat uptake, which increases the proportion of radiative forcing warming the surface. Eventually, after many centuries, surface temperature declines as the excess atmospheric carbon dioxide is taken up by the ocean and land.

## 1. Introduction

There are a wide range of climate projections for our future climate, usually based upon the surface temperature change for a cumulative carbon emission (Allen et al., 2009; Collins et al., 2013; Gillet et al., 2013; Matthews et al., 2009; Zickfeld et al., 2009). This connection provides the basis of forming budgets of the maximum carbon emission compatible with surface warming targets (Collins et al., 2013; Matthews et al., 2012; Meinshausen et al., 2009). However, it is currently unclear how the climate system responds when carbon emissions cease, either with cumulative carbon emission remaining at a maximum value or decreasing in time with carbon capture. The surface temperature change either is viewed as reaching a maximum close to when carbon emissions cease (Ricke & Caldeira, 2014) or remains nearly constant (Gillet et al., 2011; Matthews & Caldeira, 2008) or surface temperature continues to increase (Ehlert & Zickfeld, 2017; Frölicher et al., 2014; Plattner et al., 2008) over several centuries. In addition, there may be climate effects occurring after the surface temperature reaches a maximum, involving continued ocean warming and associated sea level rise (Frölicher & Joos, 2010; Gillet et al., 2011; Plattner et al., 2008). The drivers of the climate adjustment after emissions cease are unclear, although there are heuristic arguments of a partial compensation between the decline in ocean heat uptake and the reduced radiative forcing from ocean and terrestrial uptake of atmospheric CO<sub>2</sub> (Solomon et al., 2009).

In this study, we apply a theoretical relationship between global-mean surface warming and cumulative carbon emissions to understand the surface warming response on a multicentennial timescale (Goodwin et al., 2015; Williams et al., 2016). The surface warming response is considered on a timescale of many centuries after carbon emissions cease. During this period, carbon is exchanged between the atmosphere, ocean, and terrestrial biosphere; our analysis ignores sediment and weathering interactions that become significant on longer multimillennial timescales (Archer et al., 2009).

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Our theory connecting the change in global-mean surface temperature and the cumulative carbon emission is next set out (section 2). The theory is then applied to understand the surface warming response to a 1,000 year long Earth system model experiment, examining the response to carbon emissions restricted to within 100 years (Frölicher & Paynter, 2015) (section 3). The surface warming response is interpreted in terms of changes in atmospheric CO<sub>2</sub> linked to changes in the global carbon inventories and changes in the relationship between surface warming and radiative forcing. The surface warming response from the single Earth system model is compared with the response from a large ensemble of projections from an efficient Earth system model (Goodwin, 2016; Goodwin et al., 2017) (section 4). Finally, the key outcomes of the study are summarized (section 5).

## 2. Theoretical Context

If changes in radiative forcing are only driven by changes in atmospheric CO<sub>2</sub>, the global-mean change in surface temperature,  $\Delta T(t)$  (in K), may be connected via a single equation to cumulative carbon emissions (Goodwin et al., 2015), by assuming an empirical surface heat budget and combined with changes in global carbon inventories:

$$\Delta T(t) = \frac{a}{\lambda} \left( 1 - \frac{\varepsilon(t)N(t)}{R(t)} \right) \left( \frac{l_{em}(t) - \Delta l_{ter}(t) + l_{Usat}(t)}{l_B} \right). \quad (1)$$

Here the surface temperature change,  $\Delta T(t)$ , may be interpreted in terms of the product of two sets of non-dimensional time-dependent terms (each set enclosed in the brackets). The first bracket represents the fraction of radiative forcing used to warm the surface, involving the radiative forcing into the climate system,  $R(t)$  (in W m<sup>-2</sup>), the planetary heat uptake,  $N(t)$  (in W m<sup>-2</sup>), effectively from ocean heat uptake, and a nondimensional weighting of the heat uptake,  $\varepsilon(t)$  (Winton et al., 2010). The second bracket represents a carbon budget related to the change in atmospheric CO<sub>2</sub> (Goodwin et al., 2015), involving the cumulative carbon emissions,  $l_{em}(t)$ , the change in the terrestrial carbon inventory,  $\Delta l_{ter}(t)$ , and the carbon undersaturation of the global ocean,  $l_{Usat}(t)$ , and modulated by the buffered atmosphere and the ocean carbon inventory,  $l_B$  (all in PgC) (Goodwin et al., 2007). The change since the preindustrial is denoted by  $\Delta$ . The magnitude of the surface temperature change is then obtained by the product of these two sets of terms multiplied by the radiative forcing coefficient for atmospheric CO<sub>2</sub>,  $a$  (in W m<sup>-2</sup>) (Forster et al., 2013), and divided by the climate feedback parameter,  $\lambda$  (in (W m<sup>-2</sup>)K<sup>-1</sup>) (Gregory et al., 2004; Knutti & Hegerl, 2008).

If cumulative carbon emissions reach a maximum value,  $l_{em}(t_{cease})$  at a time  $t_{cease}$ , and then remain unchanged, the surface temperature change may continue to increase or decrease, based upon the evolution of the weighted ocean heat uptake,  $\varepsilon(t)N(t)$ , the radiative forcing since the preindustrial,  $R(t)$ , the ocean undersaturation of carbon,  $l_{Usat}(t)$ , and the change in the terrestrial carbon sink,  $\Delta l_{ter}(t)$ :

$$\Delta T(t) = \frac{a}{\lambda} \left( 1 - \frac{\varepsilon(t)N(t)}{R(t)} \right) \left( \frac{l_{em}(t_{cease}) - \Delta l_{ter}(t) + l_{Usat}(t)}{l_B} \right). \quad (2)$$

Surface temperature either (i) increases in time through a decline in the ratio of the weighted ocean heat uptake and radiative forcing,  $\varepsilon(t)N(t)/R(t)$ , or (ii) declines in time through the atmospheric CO<sub>2</sub> decreasing from an increase in the terrestrial carbon sink,  $\Delta l_{ter}(t)$ , or a decrease in the ocean undersaturation of carbon,  $l_{Usat}(t)$ . Thus, two sets of terms representing the thermal and carbon response of the system act in partial opposition, as argued heuristically by Solomon et al. (2009).

Eventually, the ocean approaches a long-term equilibrium with the atmosphere for both heat and carbon at a later time  $t_{equilib}$ . This long-term equilibrium state is obtained when both the ocean heat uptake,  $N(t)$ , and the global carbon undersaturation,  $l_{Usat}(t)$ , asymptote to zero in (2). For this equilibrium state (prior to any feedback from ocean sediments and weathering), the change in surface temperature is then given by the net cumulative carbon emissions to the atmosphere and ocean,  $l_{em}(t_{cease}) - \Delta l_{ter}(t_{equilib})$ , multiplied by the long-term equilibrium climate response to carbon emissions,  $a/(\lambda l_B)$  (Williams et al., 2012),

$$\Delta T(t \rightarrow t_{equilib}) \approx \frac{a}{\lambda l_B} (l_{em}(t_{cease}) - \Delta l_{ter}(t_{equilib})). \quad (3)$$

### 3. Model Analyses

Our theory is now exploited to understand the surface warming response of a 1,000 year integration of a fully coupled carbon cycle-climate Earth system model.

#### 3.1. Model Formulation and Surface Warming Simulation

We use output from a 1,000 year simulation of the global coupled carbon-climate Earth system model developed at the Geophysical Fluid Dynamics Laboratory (GFDL ESM2M) (Dunne et al., 2012; Dunne et al., 2013). The physical core of the model is an updated version of the CM2.1 (Delworth et al., 2006). The atmospheric model (AM2) has a horizontal resolution of approximately  $2^\circ$  and 24 vertical levels. The ocean model (MOM4p1) consists of 50 vertical levels and has a horizontal resolution of  $1^\circ$  or less. The ocean biogeochemical component (TOPAZv2) includes 30 tracers to represent cycles of carbon, oxygen, and the major macronutrients and iron. The land model (LM3.0) has five different vegetation pools and two soil carbon pools, and changes in these pools are simulated through phenology, natural mortality, and fire.

Following Frölicher and Paynter (2015), an idealized global warming simulation is examined where the model is forced by prescribed atmospheric  $\text{CO}_2$  increasing at an annual rate of 1% from the preindustrial 286 ppm to 745 ppm until global-mean surface air temperature increases by  $2^\circ\text{C}$  since the preindustrial (occurring in year 98 of the simulation). After that, emissions of carbon are set to zero. All other non- $\text{CO}_2$  greenhouse gases are kept at preindustrial levels.

#### 3.2. Diagnostic Approach

Based on our theory (1) (Goodwin et al., 2015; Williams et al., 2016), we diagnose the surface warming response in terms of changes in global carbon inventories and an empirical heat budget.

##### 3.2.1. Carbon Inventory Diagnostics

Our theory relies upon a global carbon inventory holding such that a cumulative carbon emission,  $I_{\text{em}}(t)$ , drives a change in the carbon inventories,

$$I_{\text{em}}(t) = \Delta I_{\text{atm}}(t) + \Delta I_{\text{ter}}(t) + \Delta I_{\text{ocean}}(t), \quad (4)$$

where  $\Delta I_{\text{atm}}(t)$ ,  $\Delta I_{\text{ter}}(t)$ , and  $\Delta I_{\text{ocean}}(t)$  are the changes in the atmospheric, terrestrial, and ocean carbon inventories since the preindustrial (all defined in PgC). Cumulative carbon emissions compatible with the prescribed atmospheric  $\text{CO}_2$  concentration are diagnosed from the sum of changes in the prescribed atmospheric  $\text{CO}_2$  inventory and the simulated time-integrated atmosphere-ocean and atmosphere-land carbon fluxes (Frölicher & Paynter, 2015; Jones et al., 2013).

The ocean carbon inventory is held as dissolved inorganic carbon,  $\Delta I_{\text{ocean}}(t) = V\Delta\text{DIC}(t)$ , where  $V$  is the global volume ( $\text{m}^3$ ) and  $\text{DIC}(t)$  is the global volume-weighted dissolved inorganic carbon concentration. A saturated component,  $\Delta C_{\text{sat}}(t)$ , of the  $\text{DIC}(t)$  is diagnosed (e.g., Lauderdale et al., 2013) based upon its potential temperature, salinity and alkalinity, and instantaneous atmospheric  $\text{CO}_2(t)$ . The carbon undersaturation of the global ocean,  $I_{\text{Usat}}$ , is then defined by  $I_{\text{Usat}}(t) = V(\Delta C_{\text{sat}}(t) - \Delta\text{DIC}(t))$  (Goodwin et al., 2015). Our theory for the carbon inventory changes uses the buffered atmosphere and ocean carbon inventory,  $I_B = I_{\text{atm}} + VC_{\text{sat}}/B$  (Goodwin et al., 2007), where  $I_{\text{atm}}$  is the atmospheric carbon inventory,  $V$  is the ocean volume,  $C_{\text{sat}}$  is the saturated dissolved inorganic carbon ( $\text{mol m}^{-3}$ ), and  $B$  is the buffer or Revelle factor (Williams & Follows, 2011); for  $\Delta I_{\text{ocean}}$ ,  $I_{\text{Usat}}$ , and  $I_B$  to be plotted in gC, there is a further multiplication of  $12 \text{ g mol}^{-1}$ .

##### 3.2.2. Radiative Forcing and Empirical Heat Budget Diagnostics

The radiative forcing,  $R(t)$  in  $\text{W m}^{-2}$ , is diagnosed from the increase in atmospheric  $\text{CO}_2$  since the preindustrial (Myhre et al., 1998),

$$R(t) = a\Delta\ln\text{CO}_2(t), \quad (5)$$

where  $a = 4.85 \text{ W m}^{-2}$  (Forster et al., 2013), and  $\Delta\ln\text{CO}_2(t) = \ln(\text{CO}_2(t)) - \ln(\text{CO}_2(t_0))$  with the preindustrial  $\text{CO}_2(t_0)$  take as 286 ppm, and a positive  $R$  represents the heat input into the climate system since the preindustrial.

The climate feedback parameter,  $\lambda$  in  $\text{W m}^{-2} \text{K}^{-1}$ , is diagnosed from a long-term empirical heat budget,

$$R(t) = \lambda\Delta T(t), \quad (6)$$

with a value of  $\lambda = 1.01 \text{ W m}^{-2} \text{ K}^{-1}$  using a least squares regression using annual-mean data from years 800 to 1000 of the 1,000 year simulation (Gregory & Forster, 2008). On shorter timescales, the radiative forcing drives a planetary heat uptake,  $N(t)$  in  $\text{W m}^{-2}$ , (Gregory et al., 2004), which is dominated by the ocean heat uptake, such that

$$R(t) = \lambda \Delta T(t) - \varepsilon(t)N(t), \quad (7)$$

where  $\varepsilon(t)$  is a nondimensional weighting of ocean heat uptake, referred to as the efficacy (Paynter & Frölicher, 2015; Winton et al., 2010, 2013). The product  $\varepsilon(t)N(t)$  is diagnosed from the mismatch of  $R(t)$  and  $\lambda \Delta T(t)$ . Alternatively, the radiative forcing may be assumed to drive a surface warming and planetary heat uptake with a time-varying climate feedback parameter,  $\alpha(t)$ , such that

$$R(t) = \alpha(t)\Delta T(t) - N(t), \quad (8)$$

where  $N(t)$  is taken from the tendency in the global ocean heat content; see Figure S1 in the supporting information for diagnostics of both these equivalent empirical heat budgets together with the time variation in  $\varepsilon(t)$  and  $\alpha(t)$ .

### 3.3. Analyses of a 1,000 Year Earth System Model Integration

The global warming response is interpreted in terms of the carbon system response, the radiative heat response, and their combined effect.

#### 3.3.1. Carbon System Response

The carbon emissions drive an increase in atmospheric  $\text{CO}_2$  from preindustrial 286 to 745 ppm (black line in Figure 1a), and an increase in the terrestrial (green line in Figure 1b) and ocean carbon inventories (blue line in Figure 1b). After emissions cease, the terrestrial and ocean systems initially both take up atmospheric  $\text{CO}_2$ , but eventually, the terrestrial uptake weakens and the ocean provides the only net carbon sink after 200 years (Figures 1b and 1c). The terrestrial carbon uptake is controlled by the land biosphere involving a competition between primary production, respiration, and disturbances from fire. By the end of the simulation, there is a small overall loss of carbon from the land biosphere due to the carbon loss from fire emissions dominating over the carbon gain due to elevated net ecosystem production.

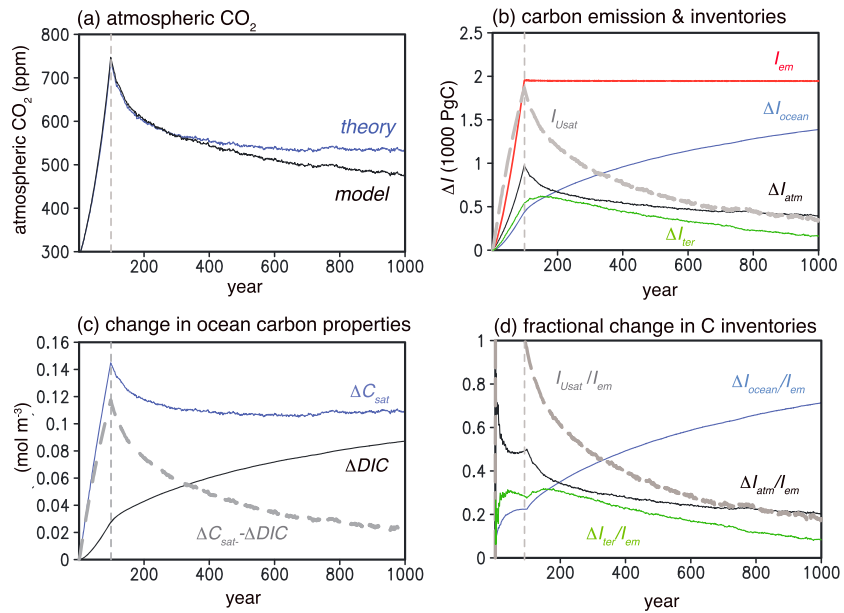
Over the 1,000 year integration, there is an overall decline in the airborne fraction,  $\Delta l_{\text{atm}}/l_{\text{em}}$  (black line in Figure 1d), an increase in the oceanborne fraction,  $\Delta l_{\text{ocean}}/l_{\text{em}}$  (blue line in Figure 1d), and an initial increase and then a later decline in the landborne fraction,  $\Delta l_{\text{ter}}/l_{\text{em}}$  (green line in Figure 1d) (Frölicher & Paynter, 2015).

The change in atmospheric  $\text{CO}_2$  may be directly connected to the net cumulative carbon emission to the combined atmosphere and ocean,  $l_{\text{em}}(t) - \Delta l_{\text{ter}}(t)$ , plus a term measuring the carbon undersaturation of the global ocean,  $l_{\text{U sat}}(t)$  (Goodwin et al., 2015), such that

$$\Delta \ln \text{CO}_2(t) = \frac{1}{l_B} (l_{\text{em}}(t) - \Delta l_{\text{ter}}(t) + l_{\text{U sat}}(t)), \quad (9)$$

where  $\Delta l_{\text{ter}}(t)$  represents the change in the terrestrial carbon inventory and  $l_B$  is the buffered atmosphere and ocean carbon inventory, measuring the available carbon in the combined atmosphere and ocean taking into account carbonate buffering, and its preindustrial value is 3,487 PgC (Goodwin et al., 2007; Williams et al., 2017). Our theoretical prediction for atmospheric  $\text{CO}_2$  from (9) agrees well with the actual model response (Figure 1a) over the first 400 years. There is a subsequent misfit due to the buffered carbon inventory,  $l_B = l_{\text{atm}} + VC_{\text{sat}}/B$ , increasing from its preindustrial value of 3,487 PgC to 4,054 PgC over the last 200 years of the record; this increase in  $l_B$  is from the effect of the rise in  $l_{\text{atm}}$  and  $C_{\text{sat}}$  exceeding the effect of the rise in  $B$  due to the large cumulative carbon emission of nearly 2,000 PgC (Goodwin et al., 2007).

The atmospheric  $\text{CO}_2$  response in (9) is primarily determined by the increase in cumulative carbon emissions,  $l_{\text{em}}(t)$ , and the change in the carbon undersaturation of the global ocean,  $l_{\text{U sat}}(t)$  (Figure 1b, red line and grey dashed line). Over the integration, the dissolved inorganic carbon,  $\Delta \text{DIC}(t)$ , increases as the ocean takes up the additional anthropogenic carbon added to the climate system (Figure 1c, black line). At the same time, the saturated carbon in the ocean,  $\Delta C_{\text{sat}}(t)$  (Figure 1c, blue line), and atmospheric  $\text{CO}_2$  increase together during the period of ongoing carbon emissions and then decrease after carbon emissions cease. The difference



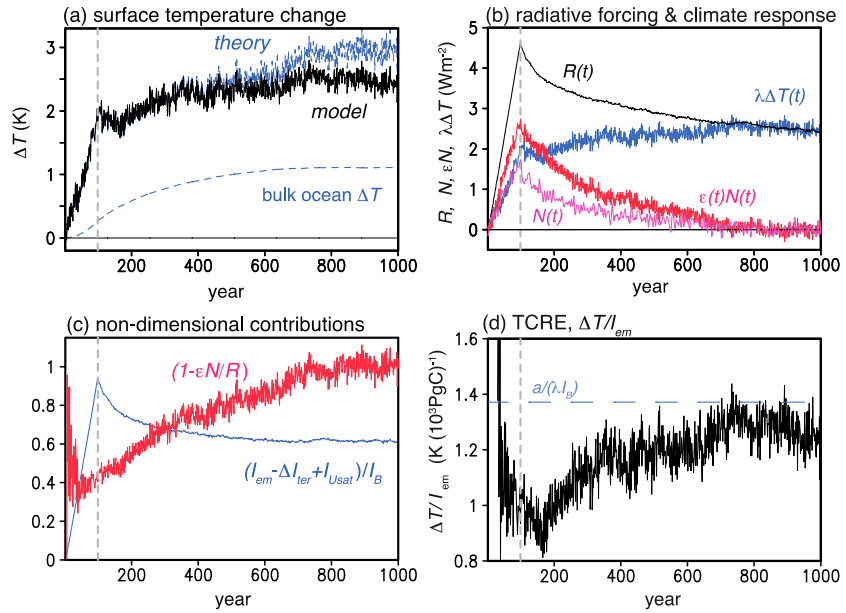
**Figure 1.** The 1,000 year experiment with the GFDL-ESM2M Earth system model: (a) change in mixing ratio for atmospheric carbon dioxide (ppm, black), peaking in year 98 and then decreasing progressively, and the prediction from our theory (ppm, blue); (b) imposed cumulative carbon emission,  $I_{em}$  (red, 1,000 PgC) with a linear increase up a 2 K warming is reached at year 98 and then no further increase (vertical dashed line denotes the cessation of emissions) and the change in global carbon inventories (in 1,000 PgC) for the atmosphere,  $\Delta I_{atm}(t)$  (black), the ocean  $\Delta I_{ocean}(t)$  (blue), and the terrestrial system  $\Delta I_{ter}(t)$  (green) together with the carbon undersaturation of the global ocean,  $I_{Usat}(t)$  (grey dashed); (c) the change in the ocean dissolved inorganic carbon,  $\Delta DIC(t)$  (black,  $\text{mol m}^{-3}$ ), and the change in the saturated dissolved inorganic carbon,  $\Delta C_{sat}(t)$  (blue,  $\text{mol m}^{-3}$ ) and their difference,  $\Delta C_{sat}(t) - \Delta DIC(t)$  (dashed grey), where the global ocean carbon undersaturation is given by  $I_{Usat}(t) = V(\Delta C_{sat}(t) - \Delta DIC(t))$  with  $V$  the global ocean volume; and (d) change in global carbon inventories divided by the cumulative carbon emissions for the atmosphere,  $\Delta I_{atm}(t)/I_{em}(t)$  (black), the ocean  $\Delta I_{ocean}(t)/I_{em}(t)$  (blue), and the terrestrial system  $\Delta I_{ter}(t)/I_{em}(t)$  (green), together with the normalized carbon undersaturation of the global ocean,  $I_{Usat}(t)/I_{em}(t)$  (dashed grey).

between the saturated carbon and dissolved inorganic carbon in the ocean,  $\Delta C_{sat}(t) - \Delta DIC(t)$ , increases rapidly during emissions and then strongly declines as the ocean takes up more carbon (Figure 1c, dashed grey line). Accordingly, the global ocean carbon undersaturation,  $I_{Usat}(t) = V(\Delta C_{sat}(t) - \Delta DIC(t))$ , declines after emissions cease (Figure 1b, dashed grey line). The normalized fraction,  $I_{Usat}(t)/I_{em}(t)$ , also strongly declines over the integration (Figure 1d, dashed grey line), which is due to the increasing effectiveness of the ocean and terrestrial system in taking up carbon on a multicentennial timescale; similar responses for  $I_{Usat}(t)/I_{em}(t)$  also occur in climate projections up to year 2100 from 11 different Earth system models (Williams et al., 2017).

### 3.3.2. Surface Warming Response

In the model experiment, carbon emissions are chosen to cease when surface temperature increases by 2 K relative to the preindustrial (this transition occurs at 98 years,  $t = t_{cease}$ ) (Frölicher & Paynter, 2015). After this period, there is still a further increase in surface temperature, rising typically by 0.6 K over the subsequent 600 years (Figure 2a).

To understand this delayed warming, we consider the empirical radiative heat budget (7). The increase in radiative forcing,  $R(t)$ , follows the rise in atmospheric  $\text{CO}_2$ , peaking at  $4.5 \text{ W m}^{-2}$  at year 98 and then declining to  $2.5 \text{ W m}^{-2}$  by year 1000 (Figure 2b, black line). This input of heat drives both a surface climate response,  $\lambda \Delta T(t)$  (blue line in Figure 2b), and a weighted planetary heat uptake,  $\varepsilon(t)N(t)$ , dominated by the ocean heat uptake (red line in Figure 2b). Initially, a larger fraction of the radiative forcing drives ocean heat uptake and a warming of the ocean interior, while only the smaller remaining fraction drives surface warming. Eventually, the ocean heat uptake declines and a larger fraction of the radiative forcing then drives a stronger surface warming. This decrease in the ocean heat uptake,  $N(t)$  (Figure 2b, purple line) is evident in the decrease in the rate of temperature rise for the bulk ocean (Figure 2a, blue dashed line).



**Figure 2.** Contributions to the surface warming in the 1,000 year experiment with the GFDL-ESM2M Earth system model: (a) the change in global-mean surface air temperature from the model,  $\Delta T(t)$  (black, K) and the predicted temperature change from our theory (1) (blue, K),  $(a/(\lambda I_B))((1 - \varepsilon(t)N(t)/R(t))(I_{em}(t) - \Delta I_{ter}(t) + I_{Usat}(t)))$ , together with the change in the bulk ocean temperature (dashed blue, K); (b) change in radiative forcing from atmospheric  $\text{CO}_2$ ,  $R(t)$  (black), surface climate response,  $\lambda \Delta T(t)$  (blue), global ocean heat uptake,  $N(t)$  (purple), and weighted ocean heat uptake,  $\varepsilon(t)N(t)$  (red), (all in  $\text{W m}^{-2}$ ); (c), the nondimensional contributions to the predicted surface temperature change,  $\Delta T(t)$ , depending on changes in the carbon inventories,  $(I_{em}(t) - \Delta I_{ter}(t) + I_{Usat}(t))/I_B$  (blue) and the global heat uptake,  $(1 - \varepsilon(t)N(t)/R(t))$  (red); and (d) the transient climate response to emissions (TCRE),  $\Delta T/I_{em}$  ( $\text{K (10}^3 \text{ PgC)}^{-1}$ ) versus year together with the equilibrium value expected from  $a/(\lambda I_B) = 1.37 \text{ K (1,000 PgC)}^{-1}$  with  $a = 4.85 \text{ W m}^{-2}$ ,  $I_B = 3,487 \text{ PgC}$ , and  $\lambda = 1.014 \text{ W m}^{-2} \text{ K}^{-1}$ .

Our theory in equation (1) provides a way to understand the surface warming response to carbon emissions (Goodwin et al., 2015); this equation is based on combining the relationship between radiative forcing and the logarithm of atmospheric  $\text{CO}_2$  (5), the empirical heat budget (7), and the relationship between the logarithm of atmospheric  $\text{CO}_2$  and the carbon inventory changes (9). The theoretical prediction of surface temperature from (9) (blue line in Figure 2a) closely matches the actual modeled surface temperature for the first 400 years (black line Figure 2a), then the theoretical prediction exceeds the model response (due to the overestimate of atmospheric  $\text{CO}_2$  over the latter part of the record; see Figure 1a).

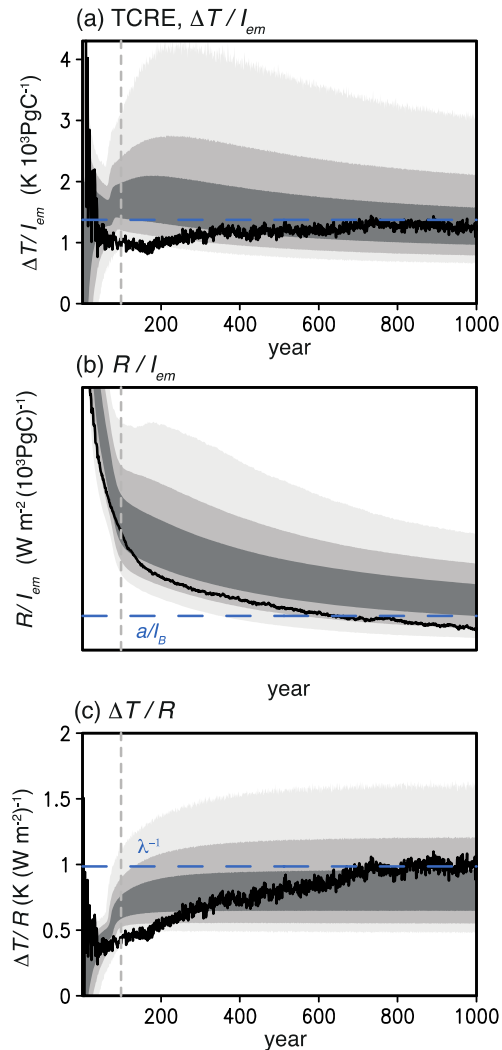
The theoretical prediction in (1) is made up of the product of two nondimensional sets of time-dependent terms: the fraction of the radiative forcing used for surface warming,  $(1 - \varepsilon(t)N(t)/R(t))$ , and the change in the logarithm of atmospheric  $\text{CO}_2$ ,  $(I_{em}(t) - \Delta I_{ter}(t) + I_{Usat}(t))/I_B$ . The contribution for atmospheric  $\text{CO}_2$  peaks when the cumulative emissions reach a maximum and then slightly declines (Figure 2c, blue line). In contrast, the fraction of radiative forcing used for surface warming initially declines, while cumulative carbon emissions increase, but then progressively increases toward 1 (Figure 2c, red line). Thus, the delay in surface warming after emission peak is due to the decline in ocean heat uptake, in accord with the model diagnostics by Frölicher et al. (2014).

### 3.3.3. Transient Climate Response to Carbon Emissions

Climate projections are often understood in terms of the transient climate response to cumulative  $\text{CO}_2$  emissions (TCRE), which is often found to be nearly constant for individual Earth system models (Allen et al., 2009; Gillet et al., 2013; Matthews et al., 2009; Zickfeld et al., 2009). The TCRE, defined here by the ratio of the surface air temperature rise,  $\Delta T(t)$ , and the cumulative carbon emission,  $I_{em}(t)$ , is equivalent to the product of the surface warming dependence on radiative forcing,  $\Delta T/R$ , and the radiative forcing dependence of cumulative carbon emissions,  $R/I_{em}$  (Williams et al., 2016):

$$\text{TCRE}(t) \equiv \frac{\Delta T(t)}{I_{em}(t)} = \left( \frac{\Delta T(t)}{R(t)} \right) \left( \frac{R(t)}{I_{em}(t)} \right), \quad (10)$$





**Figure 3.** The 1,000 year experiment with the GFDL-ESM2M Earth system model (black lines) and the efficient Earth system model (grey shading): (a) the Transient climate response to emissions (TCRE),  $\Delta T/I_{em}$  ( $K (10^3 PgC)^{-1}$ ) versus year; (b) dependence of radiative forcing on cumulative carbon emissions,  $R/I_{em}$  ( $W m^{-2} (10^3 PgC)^{-1}$ ); and (c) dependence of surface temperature on radiative forcing,  $\Delta T/R$  ( $K (W m^{-2})^{-1}$ ). The projections from a large ensemble of 10,740 efficient Earth system model configurations are denoted by grey shading (darkest, medium, and light indicate the mean  $\pm 1$ , 2 and 3 standard deviations, respectively), vertical dashed line denotes the cessation of emissions, and blue dashed lines denote the expected long-term equilibrium values for  $a/(\lambda l_B)$ ,  $a/l_B$  and  $\lambda^{-1}$  (as in Figure 2d).

which may be written from (1) (Goodwin et al., 2015) as

$$TCRE(t) \equiv \frac{\Delta T(t)}{I_{em}(t)} = \left\{ \frac{1}{\lambda} \left( 1 - \frac{\varepsilon(t)N(t)}{R(t)} \right) \right\} \left\{ \frac{a}{l_B} \left( \frac{I_{em}(t) - \Delta I_{ter}(t) + I_{Usat}(t)}{I_{em}(t)} \right) \right\}. \quad (11)$$

The TCRE varies over the 1,000 year integration: decreasing from an initial value of typically  $1.2 K (10^3 PgC)^{-1}$  to close to  $1 K (10^3 PgC)^{-1}$  when emissions cease and then increase toward  $1.3 K (10^3 PgC)^{-1}$  (Figures 2d and 3a, black line), which is close to the equilibrium value expected from  $a/(\lambda l_B) = 1.37 K (10^3 PgC)^{-1}$  (Figure 2d, blue dashed line) (Williams et al., 2012). There is significant interannual and decadal variability throughout the record.

The dependence of radiative forcing on carbon emissions,  $R/I_{em}$ , continually declines from a peak value of  $4 W m^{-2} (10^3 PgC)^{-1}$  to a value of  $1.4 W m^{-2} (10^3 PgC)^{-1}$  by year 1000 (Figure 3b). This evolution follows the response of the normalized carbon undersaturation of the global ocean,  $I_{Usat}(t)/I_{em}(t)$  (Figure 1d, grey dashed line). The carbon system is not yet at an equilibrium after year 1000 (as evident from the ratio of  $I_{Usat}(t)/I_{em}(t)$  being nonzero). The dependence of surface warming on radiative forcing,  $\Delta T/R$ , is the same as the nondimensional term  $(1 - \varepsilon(t)N(t)/R(t))$  modulated by  $\lambda^{-1}$  (Figure 3c). The surface warming dependence on radiative forcing increases in time as the weighted ocean heat uptake,  $\varepsilon(t)N(t)$ , declines in time. The thermal system is close to an equilibrium after year 1000, as evident from  $\varepsilon(t)N(t)$  being close to zero.

Hence, there are changes in the TCRE with time, but these changes are relatively small compared with the changes in the separate contributions linked to changes in the surface warming dependence on radiative forcing,  $\Delta T/R$ , and the radiative forcing dependence on cumulative carbon emissions,  $R/I_{em}$ . In comparison, there are broadly similar changes in time for climate projections to year 2100 in an ensemble of 11 Earth system models (Williams et al., 2017), although there are significant intermodel differences in the magnitudes of the TCRE and its dependences for the individual models.

## 4. Uncertainty Analysis of an Efficient Earth System Model

To provide a wider context to the Earth system model projection, we now compare the GFDL model response to a large ensemble of projections from an efficient Earth system model, designed to be consistent with observations.

### 4.1. Formulation of the Efficient Earth System Model

The Warming Acidification and Sea level Projector (WASP) model (Goodwin, 2016) is used to consider uncertainties. A large number of initial ensemble configurations ( $10^8$ ) are generated using a Monte Carlo approach, with each configuration containing a unique combination of 18 parameter values for quantities including  $a$ ,  $\lambda$ , and  $l_B$  (supporting information). Each initial configuration undergoes historical forcing from year 1765 to 2016 and is then assessed against observations for surface warming, ocean heat content change, and carbon inventory changes (Goodwin, 2016; Goodwin et al., 2017). A total of 10,470 simulations are found to be observationally consistent. These observationally consistent ensemble members are then reinitialized to a preindustrial state and forced with a 1% annual rise in atmospheric  $CO_2$  via carbon emissions into the atmosphere. As each simulation reaches a global surface temperature anomaly of  $+2.0 K$ , the carbon emissions

cease, and atmospheric  $\text{CO}_2$  is allowed to vary according to the subsequent carbon exchanges between the atmosphere, ocean, and terrestrial systems. Each separate simulation is integrated for 1,000 years.

#### 4.2. Response of the Efficient Earth System Model

The observation-consistent ensemble of model simulations reveals wide variation in the TCRE and its separate dependences for surface warming on radiative forcing and radiative forcing on carbon emissions (Figures 3a–3c, grey shading). The GFDL model response after emissions cease generally lies within the responses of the efficient model ensemble (Figures 3a–3c, black line).

Theory suggests that the TCRE,  $\Delta T/I_{\text{em}}$ , asymptotes toward the model value of  $a/(\lambda I_B)$  in (2), while  $\Delta T/R$  approaches the model value of  $\lambda^{-1}$  and  $R/I_{\text{em}}$  approaches the model value of  $a/I_B$ . In the large WASP ensemble, each simulation has a unique combination of  $a$ ,  $I_B$ , and  $\lambda$  values (Goodwin, 2016; Goodwin et al., 2017). The simulated TCRE for the GFDL model in the 1,000 year integration remains slightly below the value of  $a/(\lambda I_B)$  of  $1.37 \text{ K (1,000 PgC)}^{-1}$  (Figure 2d).

The GFDL simulation for the TCRE,  $\Delta T/I_{\text{em}}$ , lies broadly centrally within the WASP ensemble (Figure 3a) (at the 54th percentile of the WASP ensemble from year 900 to 1000). This agreement is due to the GFDL model having a relatively low value of  $a/I_B$  offsetting a high value of  $\lambda^{-1}$  compared with the WASP ensemble (Figures 3b and 3c); the GFDL model has  $a = 4.85 \text{ W m}^{-2}$  and  $I_B = 3,487 \text{ PgC}$ , placing the GFDL model value of  $a/I_B$  on the 13th percentile of the WASP ensemble, while the GFDL model has a value of  $\lambda$  of  $1.01 \text{ W m}^{-2} \text{ K}^{-1}$ , placing the GFDL model value of  $\lambda^{-1}$  on the 87th percentile of the WASP ensemble.

In the earlier stages of the integration, from years 50 to 300, the GFDL simulation of the TCRE,  $\Delta T/I_{\text{em}}$ , is lower than both the values of  $a/(\lambda I_B)$  and the response of the WASP ensemble (Figure 3a). This response is due to the relatively high value of ocean heat uptake,  $N(t)$ , in the GFDL model, which then leads to relatively low values for both the TCRE,  $\Delta T/I_{\text{em}}$ , and  $\Delta T/R$  during this period (Figures 3a and 3c).

The different transient evolution of the TCRE in the GFDL model, relative to the WASP ensemble (Figure 3a), is thus caused by differences between the transient heat uptake,  $N(t)$ , responses of the two models, since the GFDL model carbon reservoirs behavior is consistent within the WASP model ensemble (Figure 3b). This comparison highlights how the theory provides a process-based way to identify the factors determining how the TCRE evolves over time.

### 5. Conclusions

The multicentennial climate response after cessation of carbon emissions may be understood using theory connecting surface warming to cumulative carbon emissions (Goodwin et al., 2015; Williams et al., 2016), involving an empirical heat budget and global changes in carbon inventories. Our theory is compared with climate model experiments using an integration (Frölicher & Paynter, 2015) of an Earth system model (GFDL ESM2M) and a large ensemble ( $10^4$ ) of efficient Earth system models simulations (Goodwin, 2016). The surface temperature responses involves two competing contributions:

1. The surface temperature continues to increase even after a peak in atmospheric  $\text{CO}_2$  through a decline in the fraction of heat taken up by the ocean interior,  $\epsilon(t)N(t)/R(t)$ , which increases the fraction of the radiative forcing used to drive surface warming.
2. The surface temperature only decreases when there is a sufficient decrease in atmospheric  $\text{CO}_2$ , achieved through either an increase in the terrestrial carbon inventory,  $\Delta I_{\text{ter}}(t)$ , or a decline in the carbon undersaturation of the global ocean,  $I_{\text{U sat}}(t)$ .

The surface warming reaches a peak value that occurs later than the cessation of emissions due to the decline in ocean heat uptake, which initially dominates over the opposing effect of enhanced carbon uptake from the combined ocean and terrestrial system. On a millennial timescale, prior to any sediment and weathering feedback, the final surface warming is determined by the cumulative carbon emission to the combined atmosphere and ocean multiplied by a term depending upon three climate parameters,  $a/(\lambda I_B)$  (Williams et al., 2012).

This study reveals how the theory may be used to understand the response of Earth system models. Within our theory, these intermodel differences are represented by differences in ocean heat and carbon uptake,



terrestrial carbon cycling, the climate sensitivity, and the efficacy, representing the nondimensional weighting of heat uptake. Time dependence in the feedbacks (Knutti & Rugenstein, 2015) is likely to lead to the climate sensitivity or efficacy evolving in time, such as in how climate sensitivity may vary through regional feedbacks (Armour et al., 2013) and in how the efficacy may alter with changes in ocean circulation (Rose & Rayborn, 2016; Winton et al., 2013).

The delayed warming after carbon emissions cease investigated here provides an additional challenge as to how policymakers should approach restricting warming to predefined targets, such as the Paris Agreement 1.5 and 2.0 K targets above preindustrial (Ehlert & Zickfeld, 2017; Matthews et al., 2012; Meinshausen et al., 2009). Policy makers are usually focusing on the risks of climate hazards over the next century, but there are longer-term climate threats (Friedlinstein et al., 2011; Frölicher et al., 2014) from continued surface warming after carbon emissions cease. The time-dependent drivers of these climate threats after emissions cease are then primarily related to the long-term climate impacts of how the ocean takes up heat and carbon.

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## Drivers of continued surface warming after cessation of carbon emissions

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#### Text S1.

The empirical heat budget may be written in two equivalent ways:

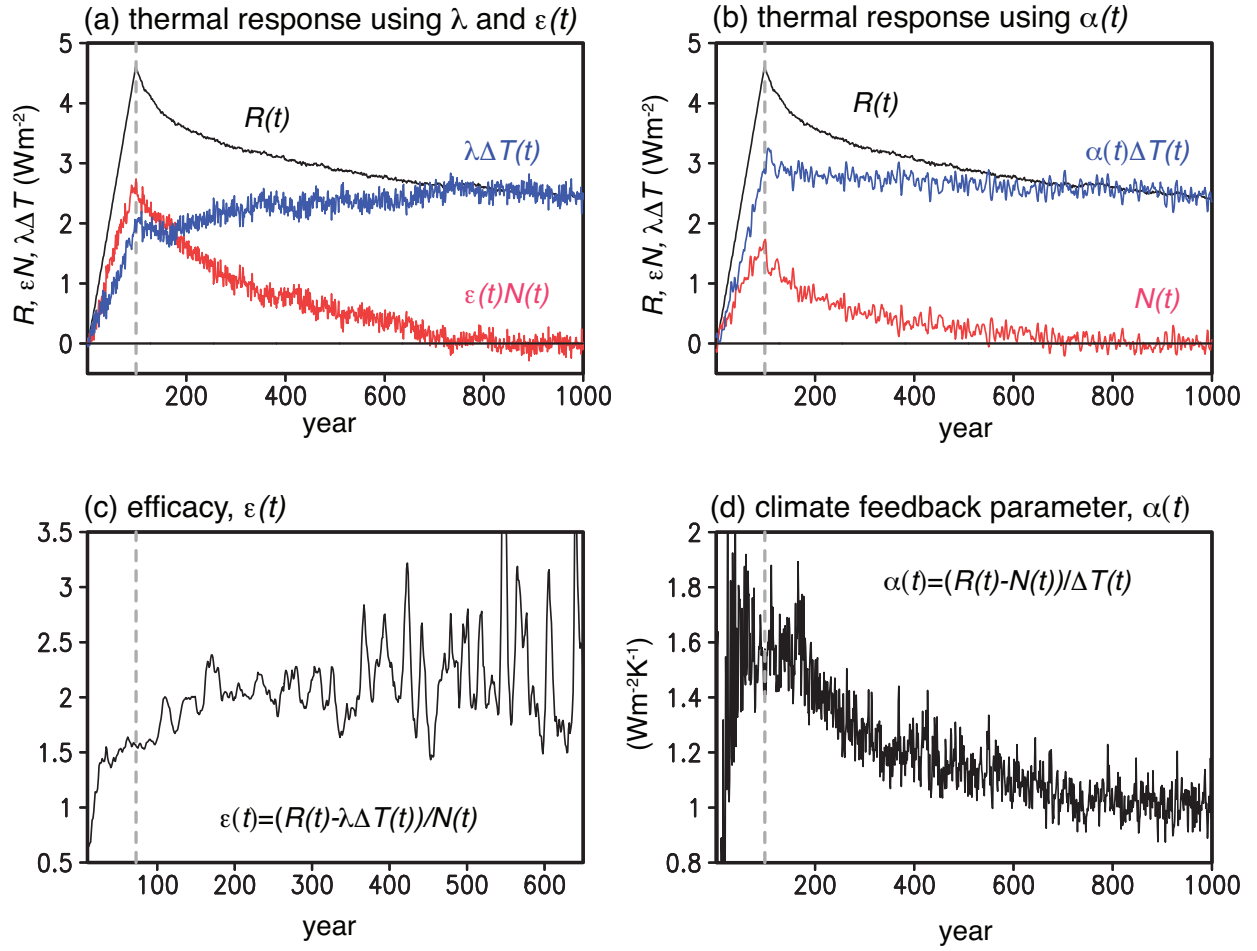
$$R(t) = \lambda \Delta T(t) + \varepsilon(t)N(t) = \alpha(t)\Delta T(t) + N(t),$$

where  $\lambda$  is the equilibrium climate feedback parameter ( $\text{Wm}^{-2}\text{K}^{-1}$ ) and is chosen to be fixed in time,  $\alpha(t)$  is a climate feedback parameter ( $\text{Wm}^{-2}\text{K}^{-1}$ ) allowed to vary in time,  $\varepsilon(t)$  is a non-dimensional weighting of planetary heat uptake allowed to vary in time,  $R(t)$  is the radiative forcing since the preindustrial ( $\text{Wm}^{-2}$ ),  $\Delta T(t)$  is the global-mean surface air temperature since the preindustrial (K), and  $N$  is the planetary heat uptake ( $\text{Wm}^{-2}$ ) that is effectively measured from ocean heat uptake.

The radiative forcing from atmospheric  $\text{CO}_2$ ,  $R(t)$ , drives a surface climate response and planetary heat uptake, which may be represented by  $\lambda \Delta T(t)$  plus the weighted ocean heat uptake,  $\varepsilon(t)N(t)$  (Fig. S1a,c), or by  $\alpha(t)\Delta T(t)$  plus the ocean heat uptake,  $N(t)$  (Fig. S1b,d). If the latter form is adopted, then the relationship between surface temperature change and carbon emissions given in equation (1) may be rewritten as

$$\Delta T(t) = \frac{a}{\alpha(t)I_B} \left( 1 - \frac{N(t)}{R(t)} \right) (I_{em}(t) - \Delta I_{ter}(t) + I_{Usat}(t))$$

Time dependence in the feedbacks [Knutti and Rugenstein, 2017] are likely to lead to the climate feedback parameter,  $\alpha(t)$ , and efficacy,  $\varepsilon(t)$ , evolving in time; such as in how the climate feedback parameter varies through regional feedbacks [Armour *et al.*, 2013] or in how the efficacy alters with changes in ocean circulation [Winton *et al.*, 2013; Rose and Rayborn, 2016].



**Supplementary Figure 1.** Empirical heat budget over time: (a) the radiative forcing from atmospheric CO<sub>2</sub>,  $R(t)$  (black), balances the surface climate response,  $\lambda\Delta T(t)$  (blue), plus the weighted ocean heat uptake,  $\varepsilon(t)N(t)$  (red), and (b) the radiative forcing from atmospheric CO<sub>2</sub>,  $R(t)$  (black), balances the surface climate response,  $\alpha(t)\Delta T(t)$  (blue), plus the global ocean heat uptake,  $N(t)$  (red); all in  $\text{Wm}^{-2}$ . In addition, in (c) the time-dependent non-dimensional weighting of the heat uptake, the efficacy,  $\varepsilon(t)$ , over the first 600 years when the product  $\varepsilon(t)N(t)$  is non zero, and in (d) the time-variation of the time-dependent climate feedback parameter,  $\alpha(t)$  ( $\text{Wm}^{-2}\text{K}^{-1}$ ) over the entire record.

## Text S2.

### The WASP Earth system model ensemble

The Warming Acidification and Sea level Projector (WASP) is used to generate an initial ensemble of  $10^8$  simulations in a Monte Carlo approach following *Goodwin* [2016] and *Goodwin et al.* [2017]. A total of 18 model parameters are varied between simulations [*Goodwin*, 2016; *Goodwin et al.*, 2017]:  $a$ ;  $\lambda$ ;  $I_B$ ; 5 ocean tracer uptake timescales [*Goodwin*, 2017]; 2 equilibrium warming ratios (the ratio of warming between global air temperatures and sea surface temperatures, and the ratio of warming between sea surface temperatures and sub-surface ocean temperatures); the fraction of anthropogenic heat content increase that enters the ocean; 2 uncertainty scaling parameters related to uncertainty in radiative forcing from agents other than  $\text{CO}_2$  (one for other greenhouse gasses and one for aerosols); and 2 non-dimensional efficacies relating the relative warming of processes to that incurred by  $\text{CO}_2$  radiative forcing ( $\epsilon$  representing the relative warming effect of ocean heat uptake and  $\epsilon_{aero}$  representing the relative warming effect of radiative forcing from aerosols).

Each of the  $10^8$  simulations are forced with historic  $\text{CO}_2$  and radiative forcing from other agents, from 1765 up to 2016, and assessed against reconstructions of historic changes to the real climate system (for surface warming, ocean heat content change and ocean and terrestrial carbon uptake) derived from observations. Following the methodology of *Goodwin* [2016] and *Goodwin et al.* [2017], observational-consistent ranges for each constraint (e.g. surface warming) are based on the 90% confidence ranges for the quantity in the real climate system. A WASP simulation is accepted as observationally consistent if the simulated observables (based on surface warming, ocean heat content changes and ocean and terrestrial carbon uptake) fit either within the ranges of all observational constraints, or fit within all but one constraint and are less than 50% outside the range of the final constraint. This assessment allows the tails of the distributions for surface warming and ocean heat content to be included within the observation-consistent WASP ensemble.

After the observational tests, some 10,740 simulations remain, or 0.01% of the initial Monte Carlo simulations. These observationally-consistent ensemble members are then separately re-initialised and forced with the 1% annual rise in  $\text{CO}_2$  experiments to compare to the GFDL model.

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