

The role of non-Newtonian fluids in microswimmer propulsion

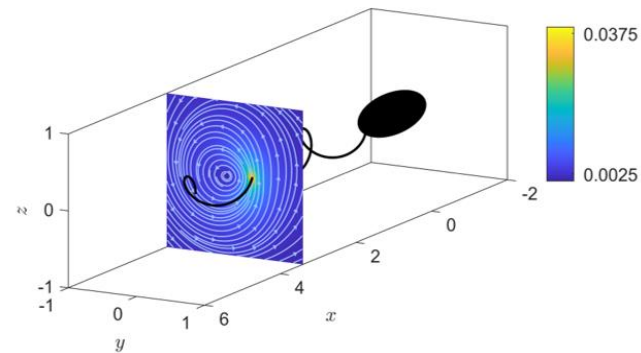
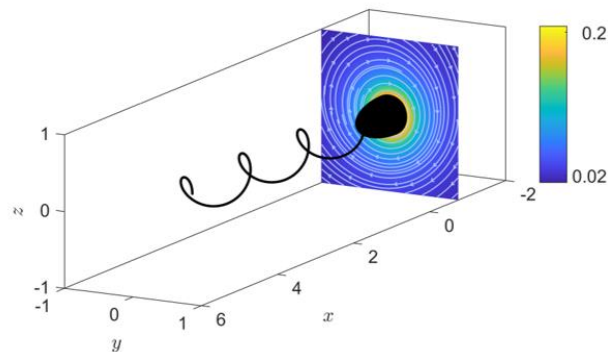
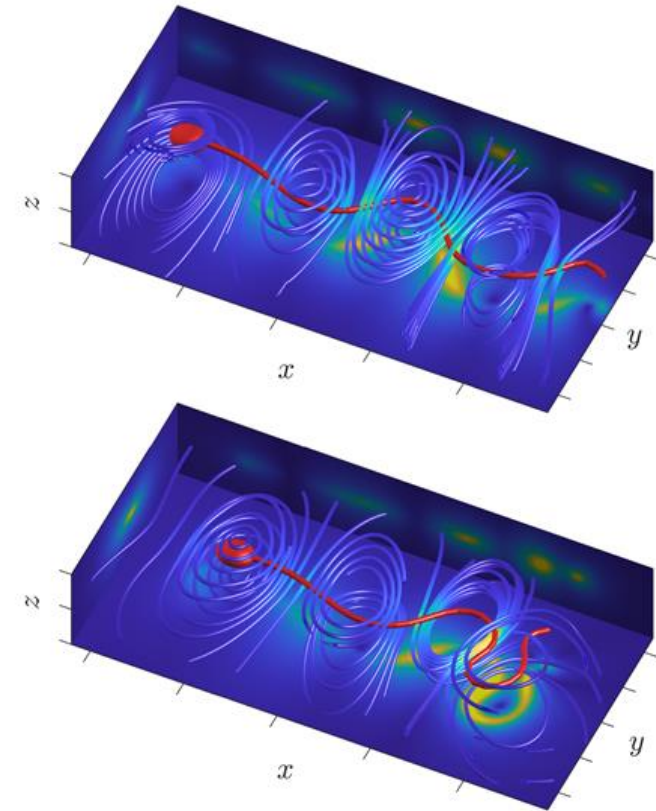
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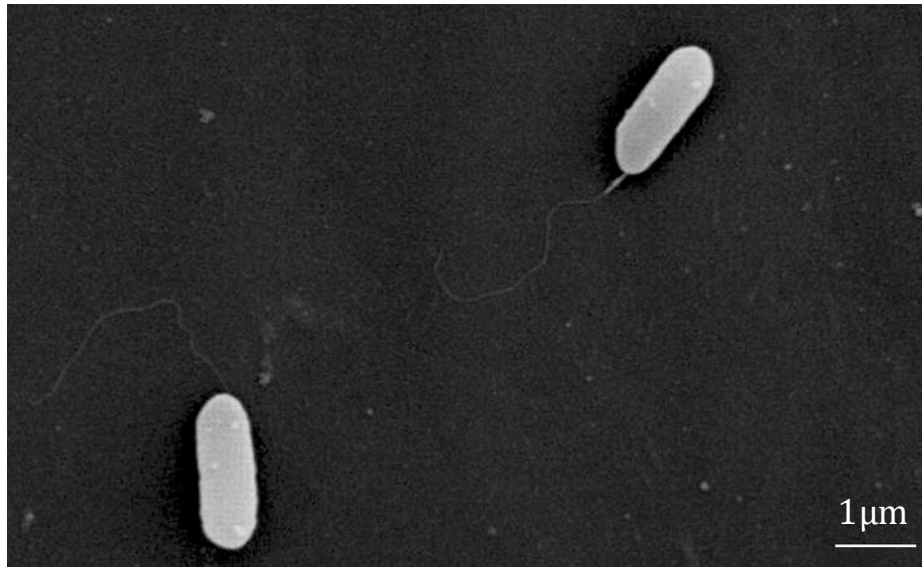
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Swimming in non-Newtonian fluids

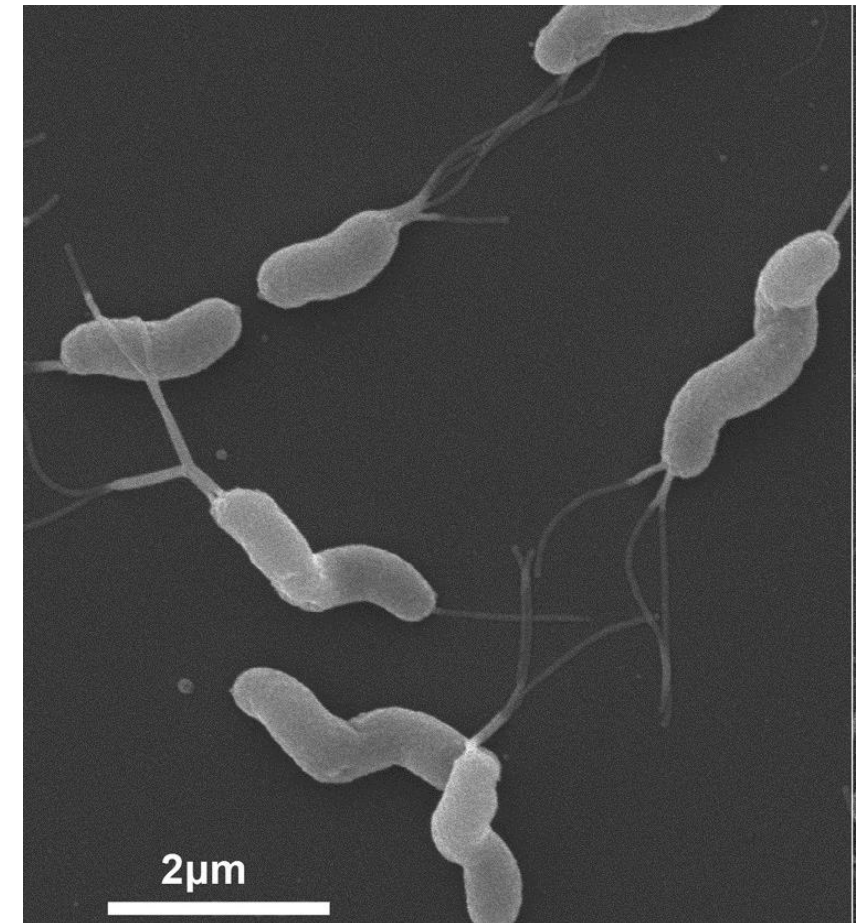


1. *Pseudomonas aeruginosa* progressing through the mucus-filled respiratory system.

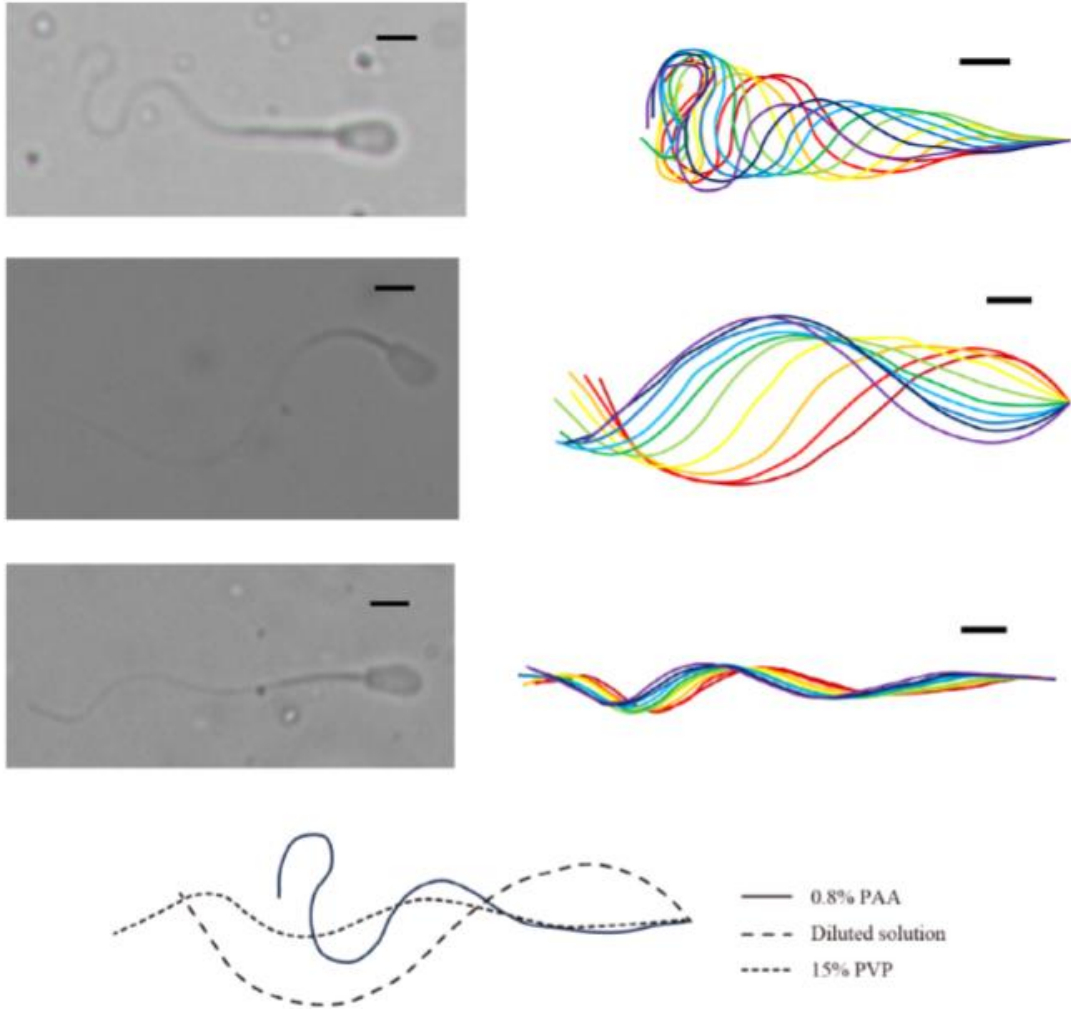
3. Human spermatozoa progressing through a cervical mucus analogue.



2. *Helicobacter pylori* moving through the mucus layer covering the stomach.

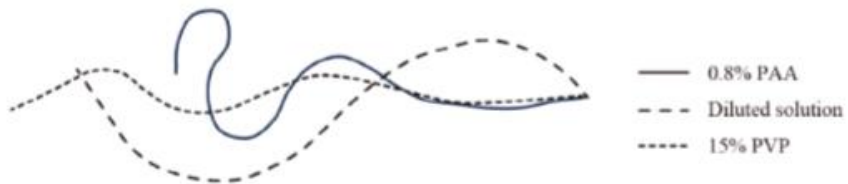
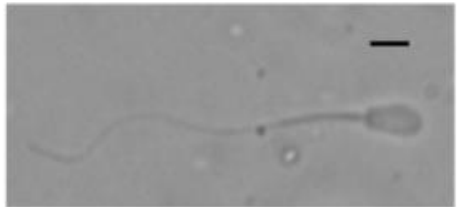
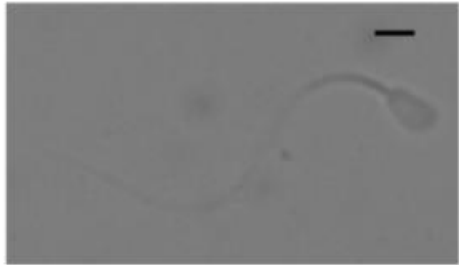


Defining the fluid-structure problem



Hyakutake, Sato & Sugita. *J. Biomech.* 2019.

Defining the fluid-structure problem



Hyakutake, Sato & Sugita. *J. Biomech.* 2019.

$$-\nabla p + \nabla \cdot \boldsymbol{\tau} = \mathbf{0}$$

$$\nabla \cdot \mathbf{u} = 0$$

$Re = \frac{\rho UL}{\mu} \rightarrow 0$, steady,
 incompressible

Newtonian

$$\boldsymbol{\tau} = 2\mu \mathbf{D}(\mathbf{u})$$

Viscoelastic

$$\boldsymbol{\tau} + \lambda_1 \overset{\nabla}{\boldsymbol{\tau}} = 2\mu_0 [\mathbf{D}(\mathbf{u}) + \lambda_2 \overset{\nabla}{\mathbf{D}}(\mathbf{u})]$$

Shear-dependent viscosity

$$\boldsymbol{\tau} = 2\mu(\dot{\gamma}) \mathbf{D}(\mathbf{u})$$

The hybrid approach to non-linear swimming

The boundary value problem

$$\begin{aligned}\nabla \cdot [2\mu(\dot{\gamma}(\mathbf{u}))\mathbf{D}(\mathbf{u})] - \nabla p + \mathbf{F} &= \mathbf{0} && \text{in } \Omega \\ \nabla \cdot \mathbf{u} &= 0 && \text{in } \Omega \\ \mathbf{u} &= \mathbf{0} && \text{on } \partial\Omega\end{aligned}$$

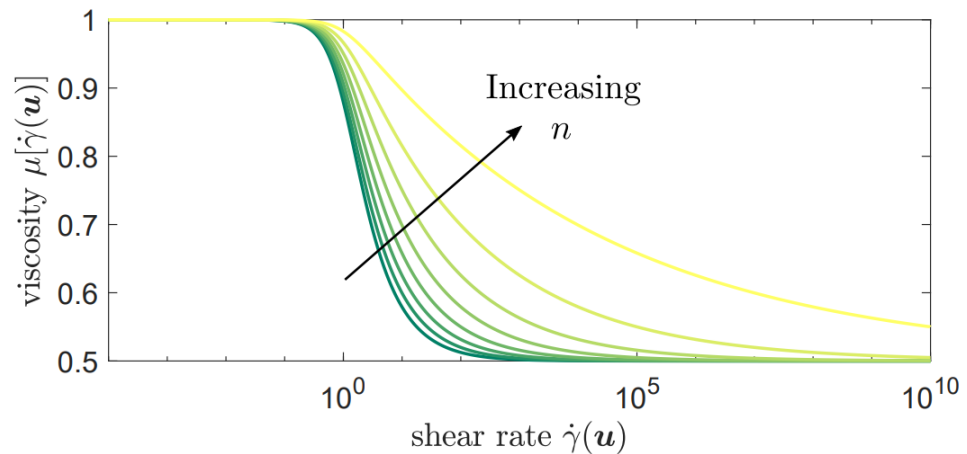
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Constitutive law (Bird-Carreau-Cross)

$$\mu(\dot{\gamma}) = \mu_{\infty} + [\mu_0 - \mu_{\infty}][1 + [\lambda\dot{\gamma}]^2]^{\frac{n-1}{2}}$$

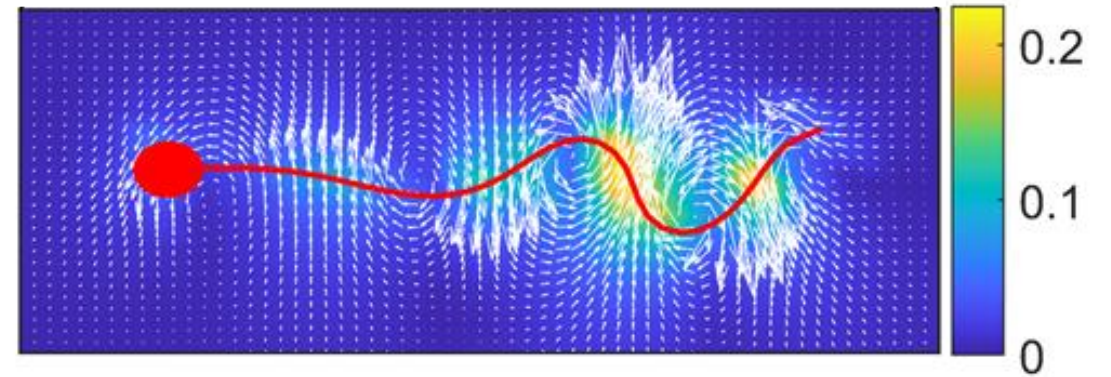


The hybrid approach to non-linear swimming

The boundary value problem

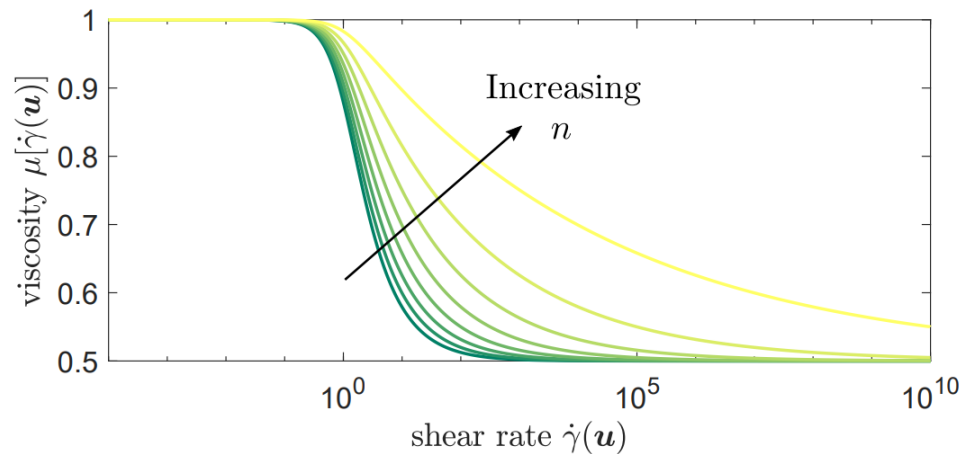
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Fluid flow around a swimmer varies rapidly

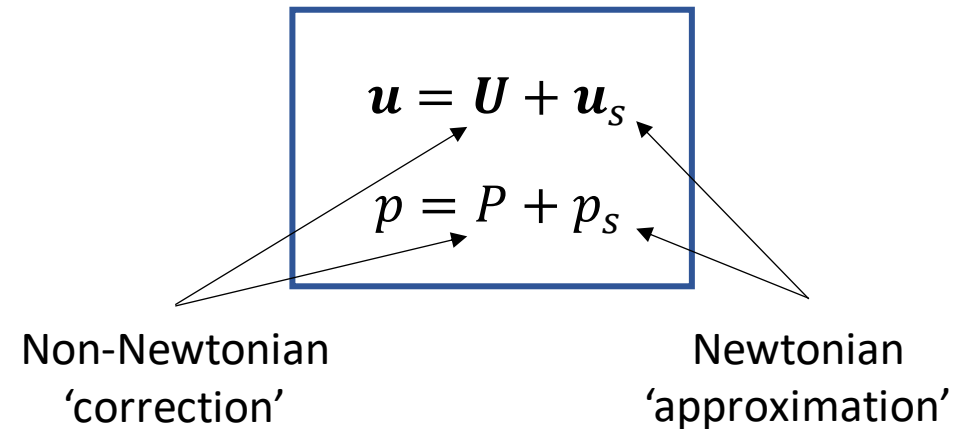


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The hybrid method



The hybrid approach to non-linear swimming: example

The boundary value problem

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Newtonian flow approximation

If $\mathbf{F} = \mathbf{f}\delta(\mathbf{x} - \mathbf{y})$, (when $\mu(\dot{\gamma}) := \mu_N$)

$$\mathbf{u}_s = S_{ij}(\mathbf{x}, \mathbf{y})\mathbf{f}_j = \frac{1}{8\pi\mu_N} \left[\frac{\delta_{ij}}{r} + \frac{r_i r_j}{r^3} \right] \mathbf{f}_j,$$

$$p_s = P_j(\mathbf{x}, \mathbf{y})\mathbf{f}_j = \frac{1}{4\pi} \left[\frac{r_j}{r^3} \right] \mathbf{f}_j,$$

where $r_i = x_i - y_i$, $r = |\mathbf{x} - \mathbf{y}|$.

The hybrid approach to non-linear swimming: example

The boundary value problem

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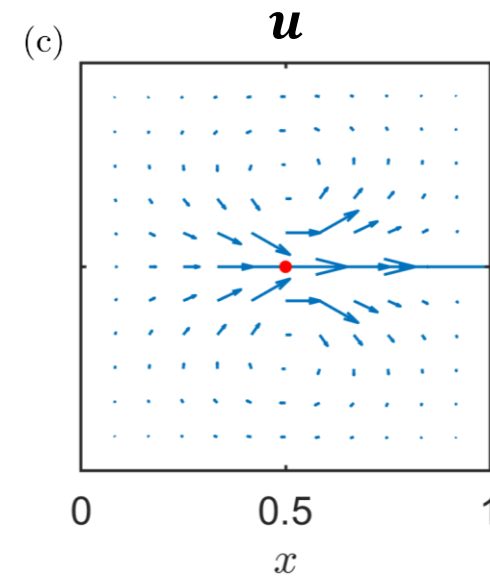
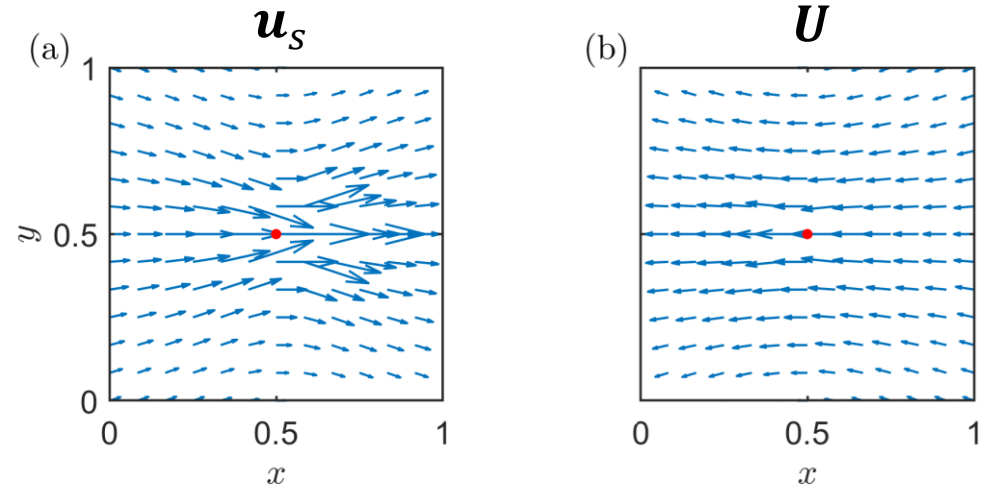
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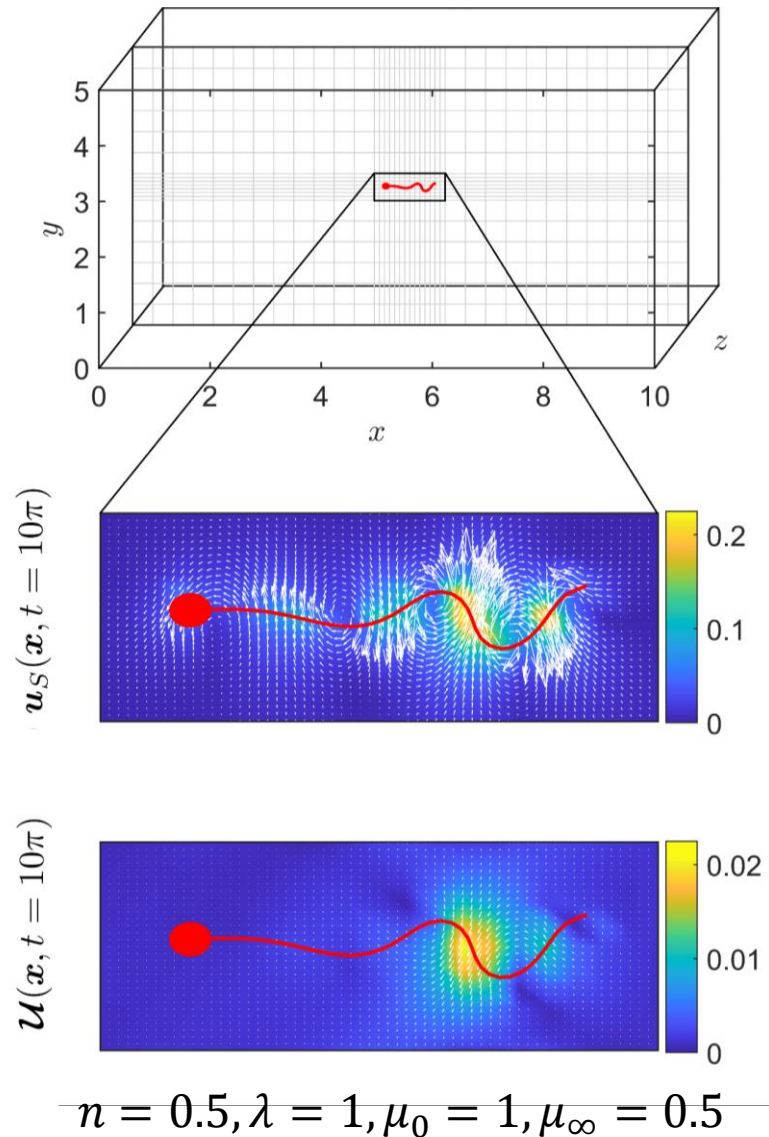
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$$n = 0.5, \lambda = 1, \mu_0 = 1, \mu_\infty = 0.5$$

The effect of shear-thinning rheology on sperm propulsion

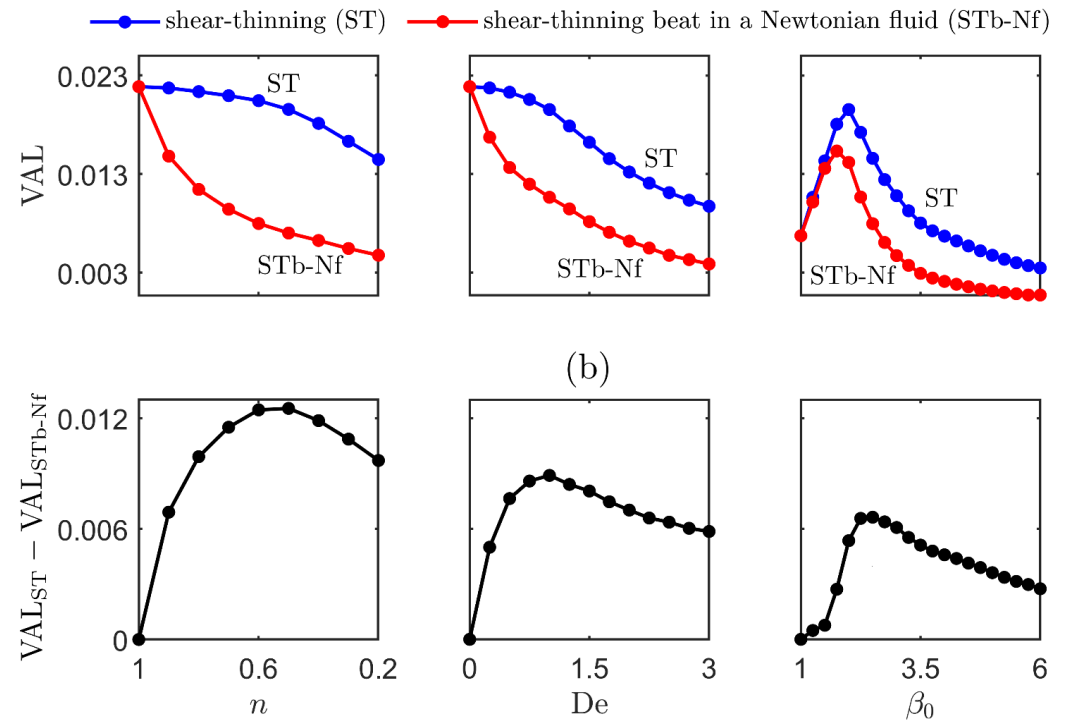


$$\mu(\dot{\gamma}) = 1 + [\beta_0 - 1][1 + [\text{De } \dot{\gamma}]^2]^{\frac{n-1}{2}}$$

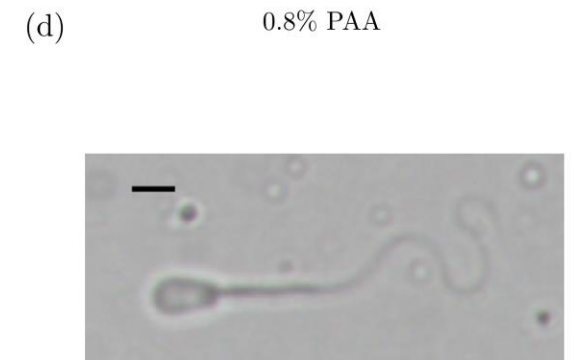
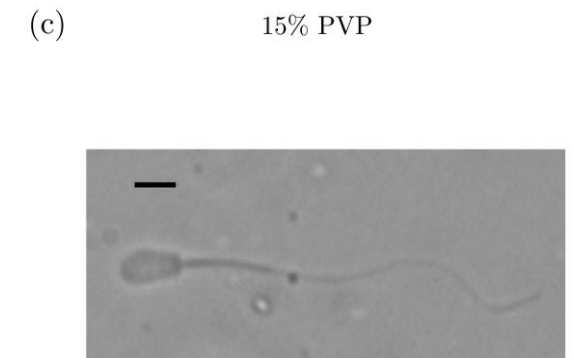
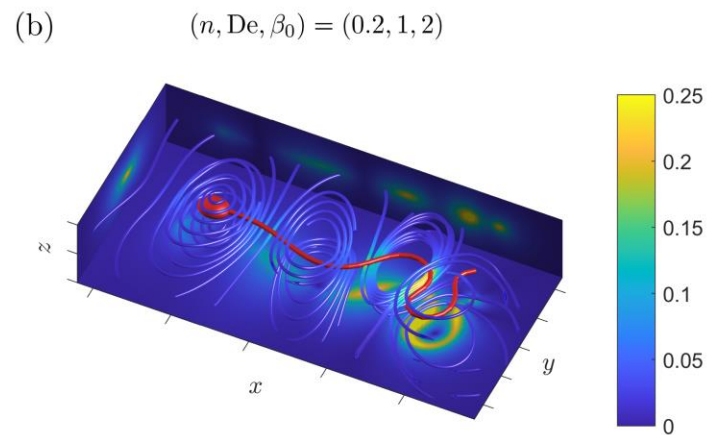
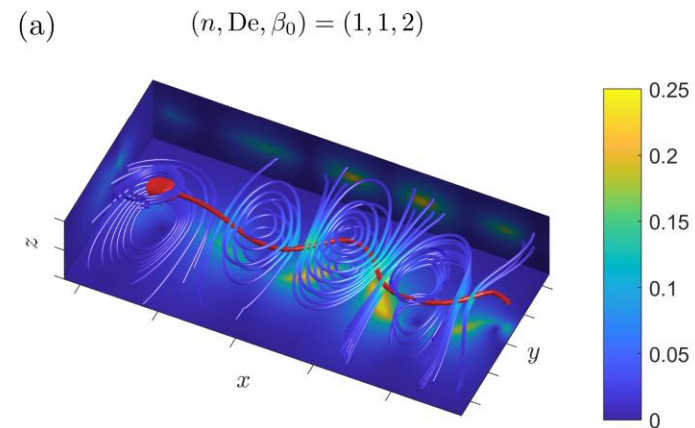
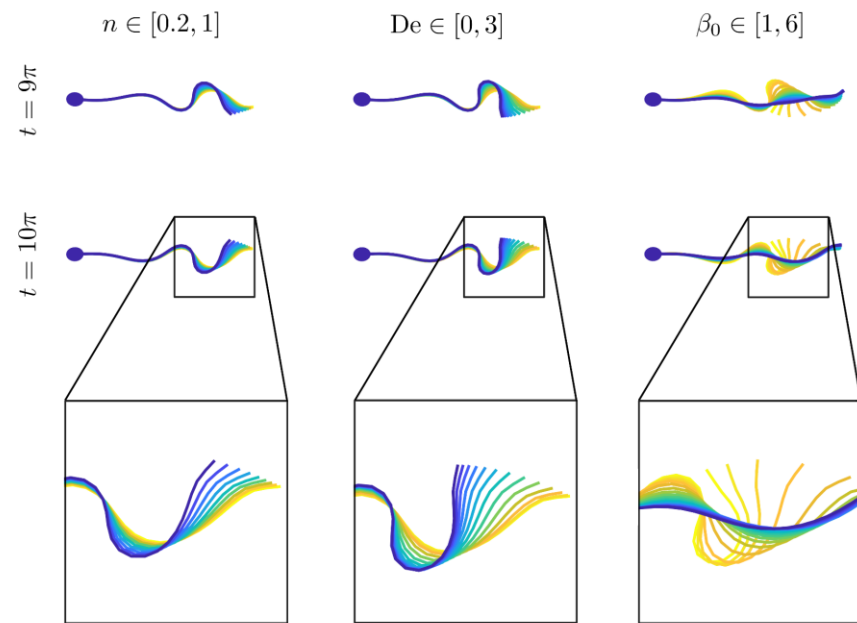
$$\text{De} = \lambda\omega$$

$$\beta_0 = \frac{\mu_0}{\mu_\infty}$$

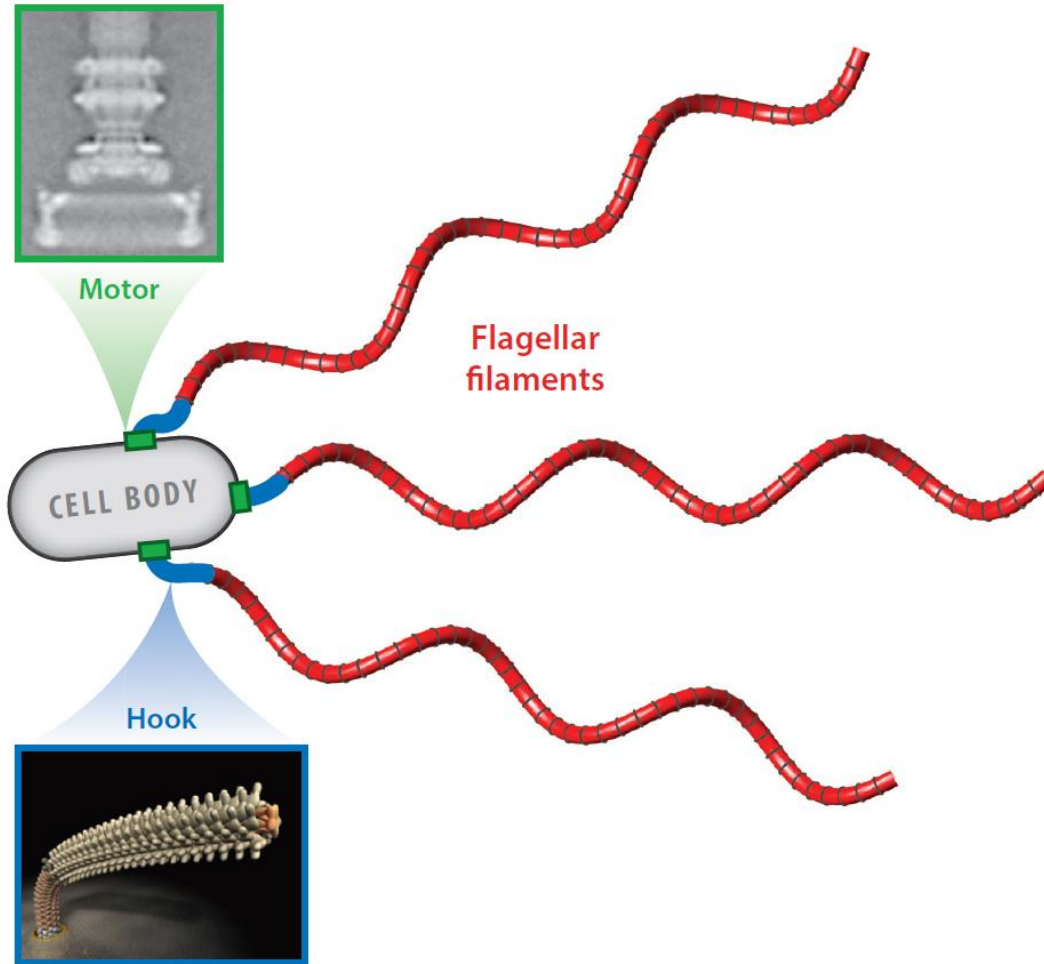
(a)



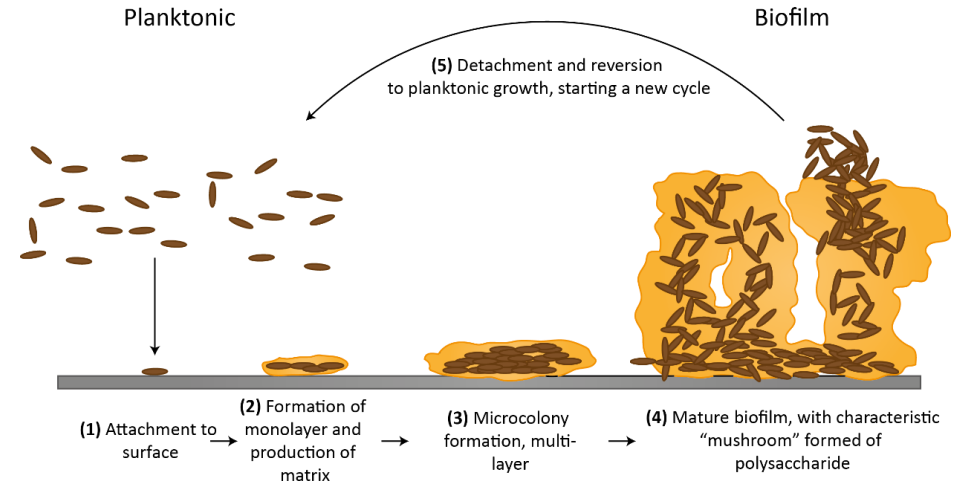
The effect of shear-thinning rheology on sperm propulsion



Bacterial locomotion and biofilm formation in non-Newtonian fluids



E. Lauga. *Annu. Rev. Fluid Mech.* 2016.



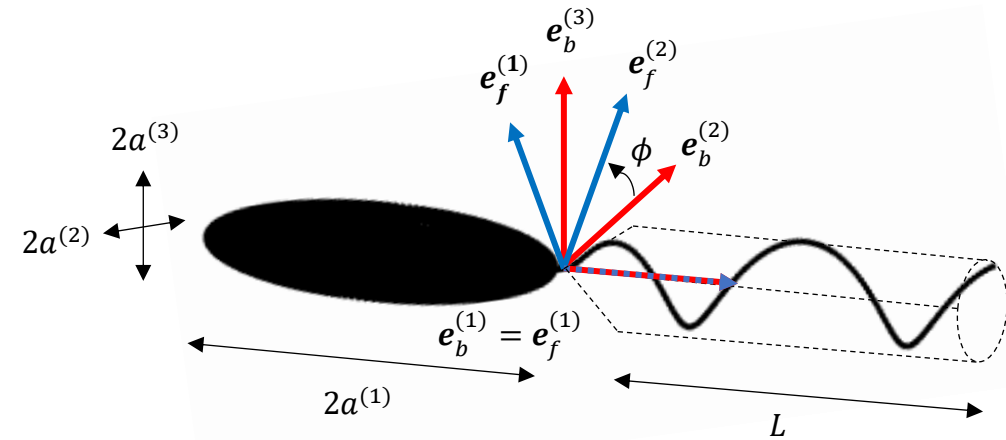
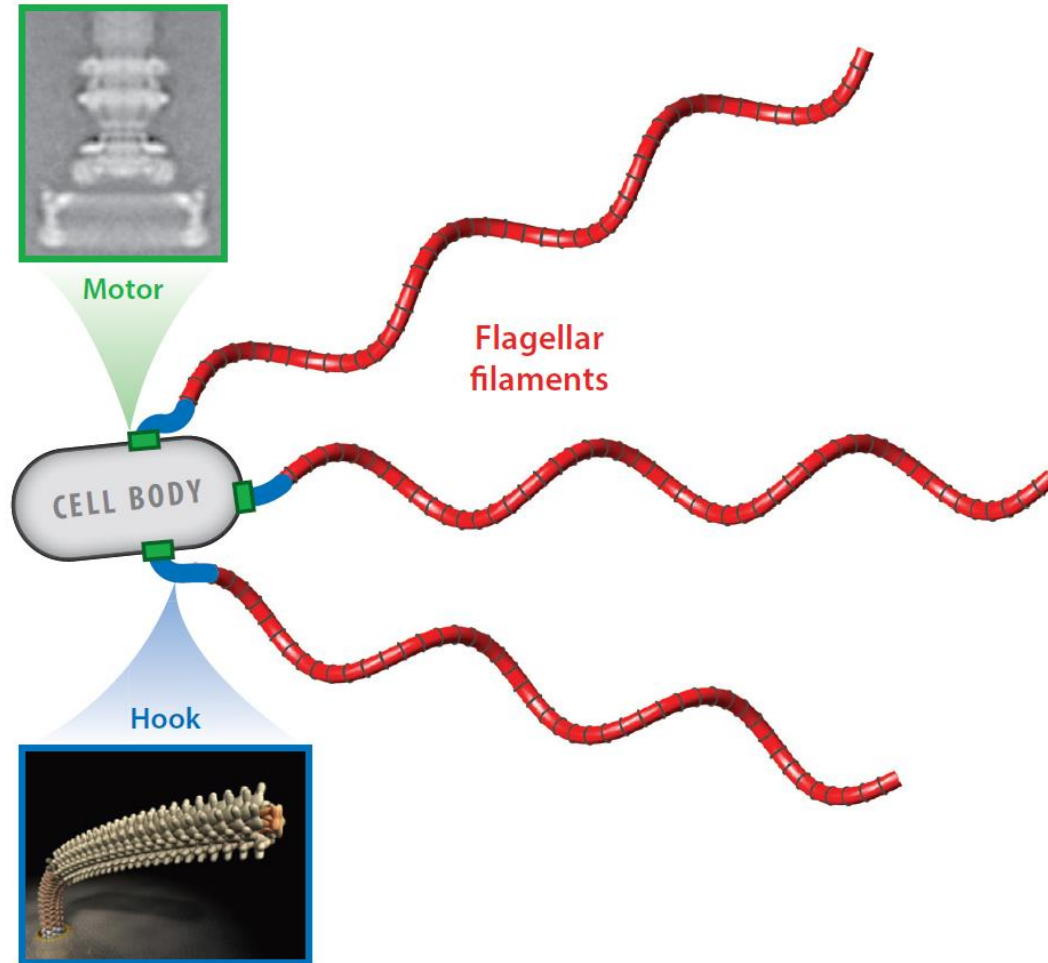
Vasudevan. *J. Microbiol. Exp.* 2014.

$$\begin{aligned}
 \mathbf{u} &= \mathbf{U} + \mathbf{u}_s \\
 p &= P + p_s
 \end{aligned}$$

Alter: viscoelastic + shear-thinning effects

Alter: bacterium model with rigid flagellum

The bacterium model



Motion is driven through a prescribed relative rotation rate between the body and flagellum

$$\omega = \frac{d\phi}{dt}.$$

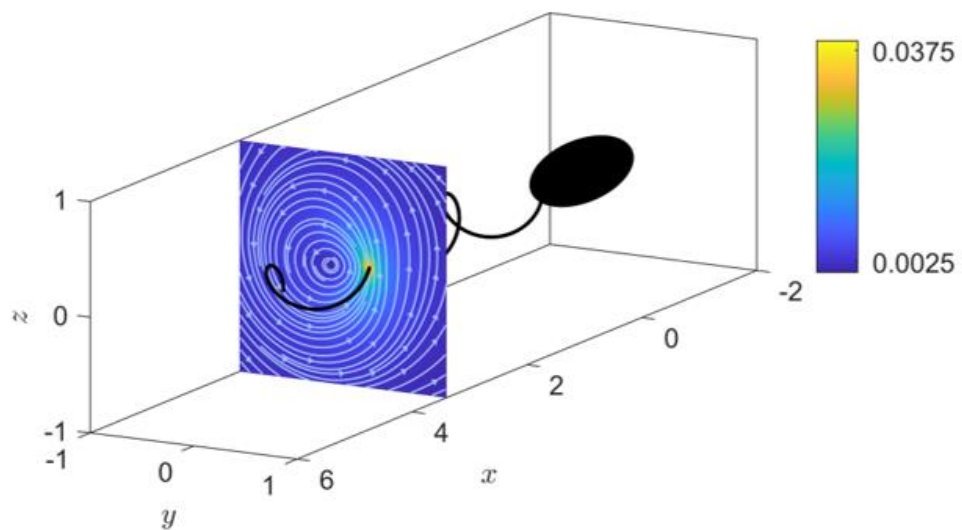
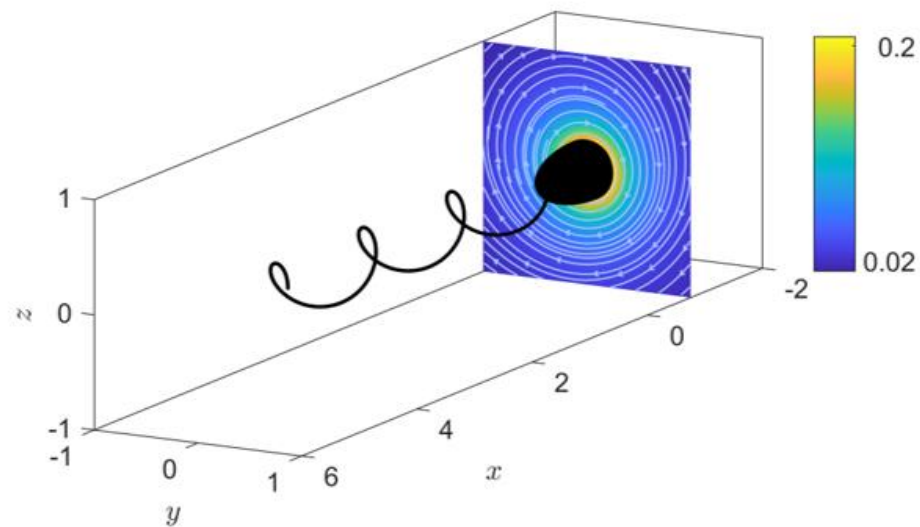
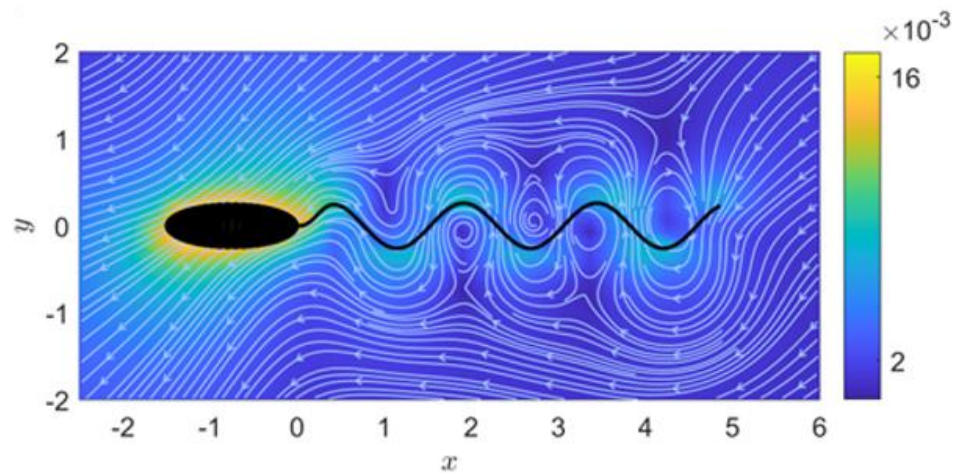
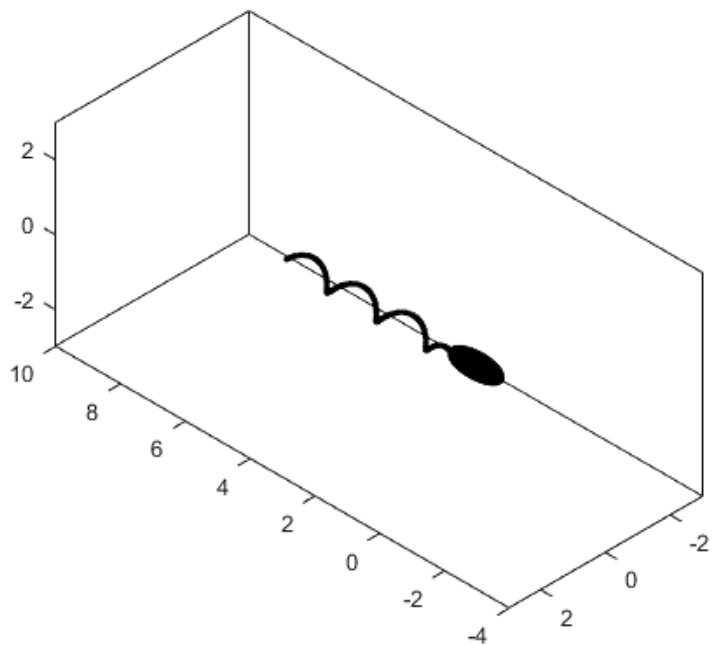
The surface velocity of points on the body and flagellum are described by

$$u(\mathbf{Y}) = \mathbf{V} + \boldsymbol{\Omega}_b \times [\mathbf{Y} - \mathbf{X}_0],$$

$$u(\mathbf{X}) = \mathbf{V} + [\boldsymbol{\Omega}_b + \boldsymbol{\Omega}_m] \times [\mathbf{X} - \mathbf{X}_0].$$

$$(\boldsymbol{\Omega}_m = [\omega, 0, 0]^T)$$

The bacterium model

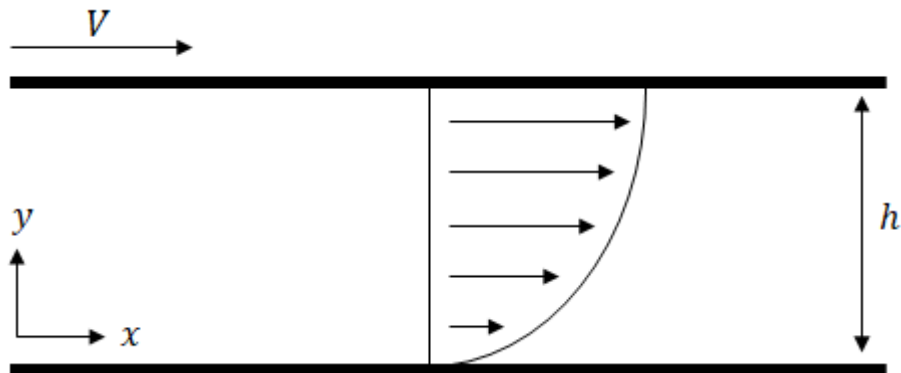


Couette-Poiseuille flow of a Giesekus fluid

Giesekus model

$$\boldsymbol{\tau} + \frac{\alpha\lambda}{\eta}(\boldsymbol{\tau} \cdot \boldsymbol{\tau}) + \lambda \overset{\nabla}{\boldsymbol{\tau}} = 2\eta \mathbf{D}$$

The parameters λ and η are the fluid relaxation time and viscosity respectively. The parameter $0 < \alpha < 1$ is the dimensionless mobility.

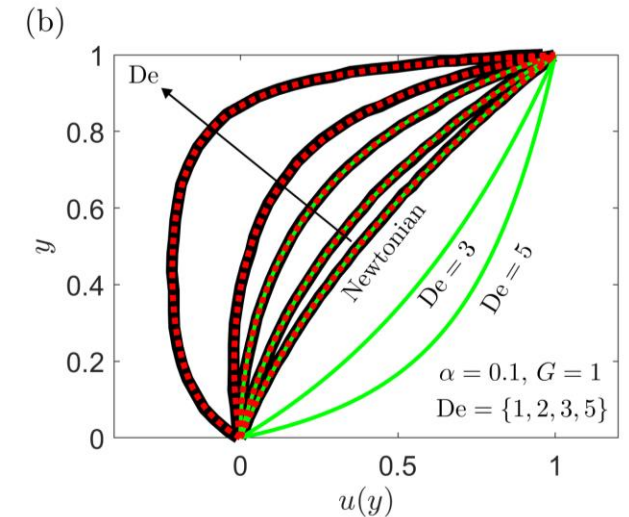
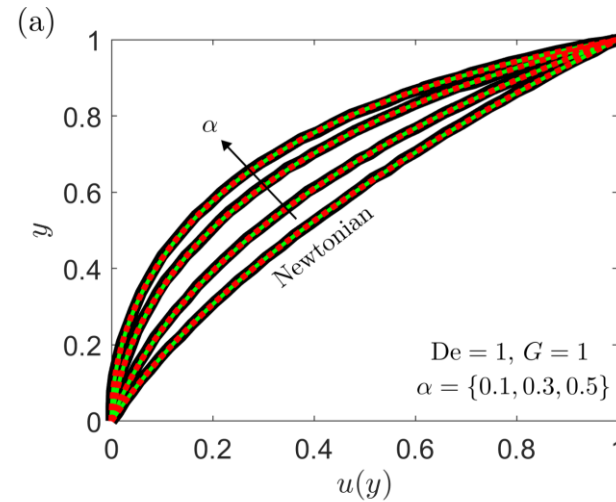
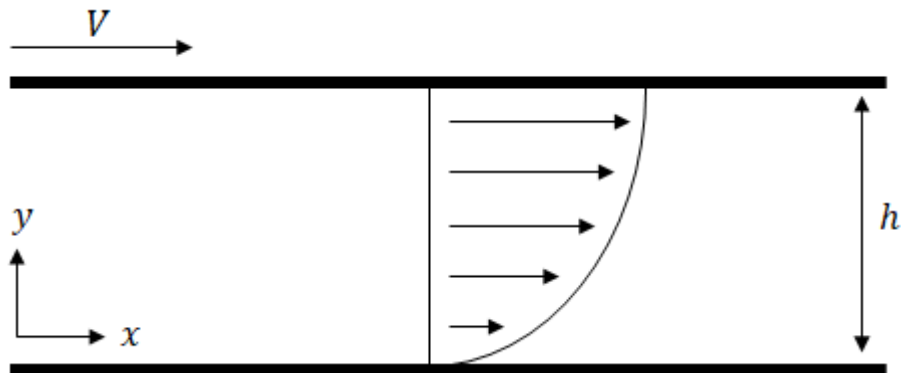


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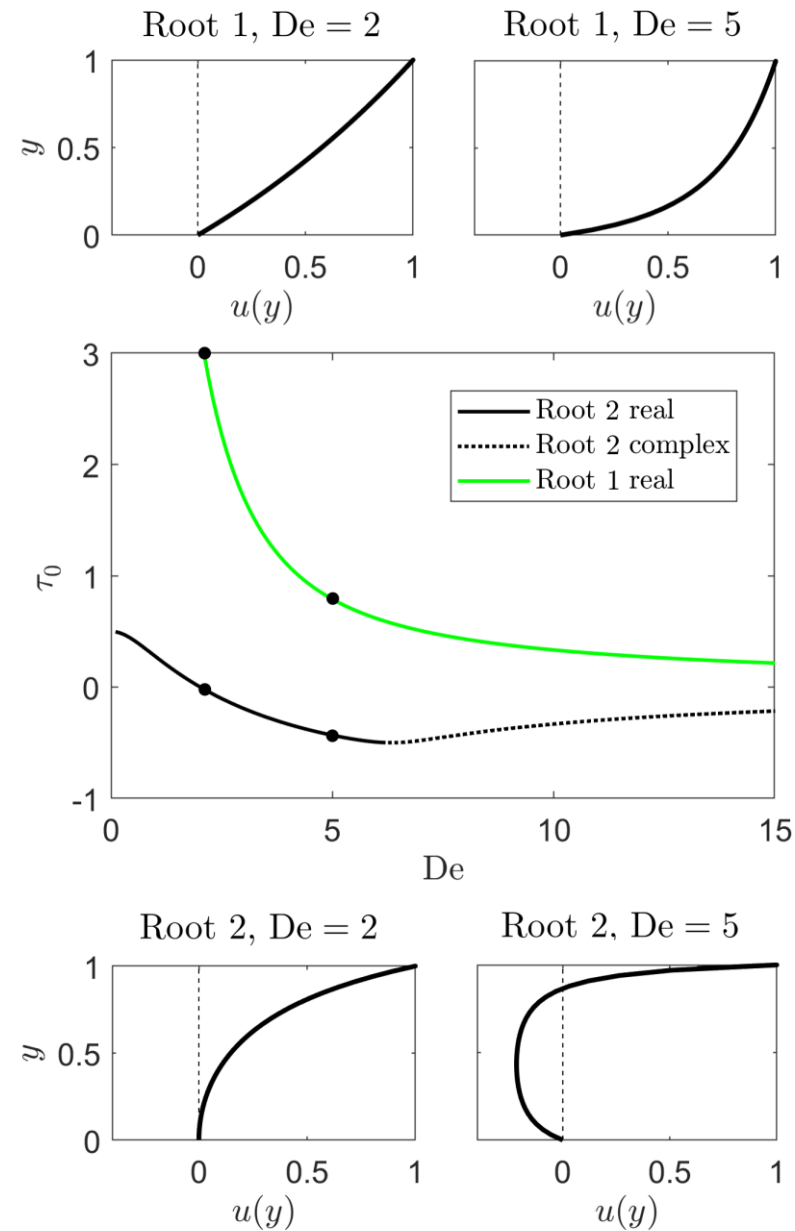
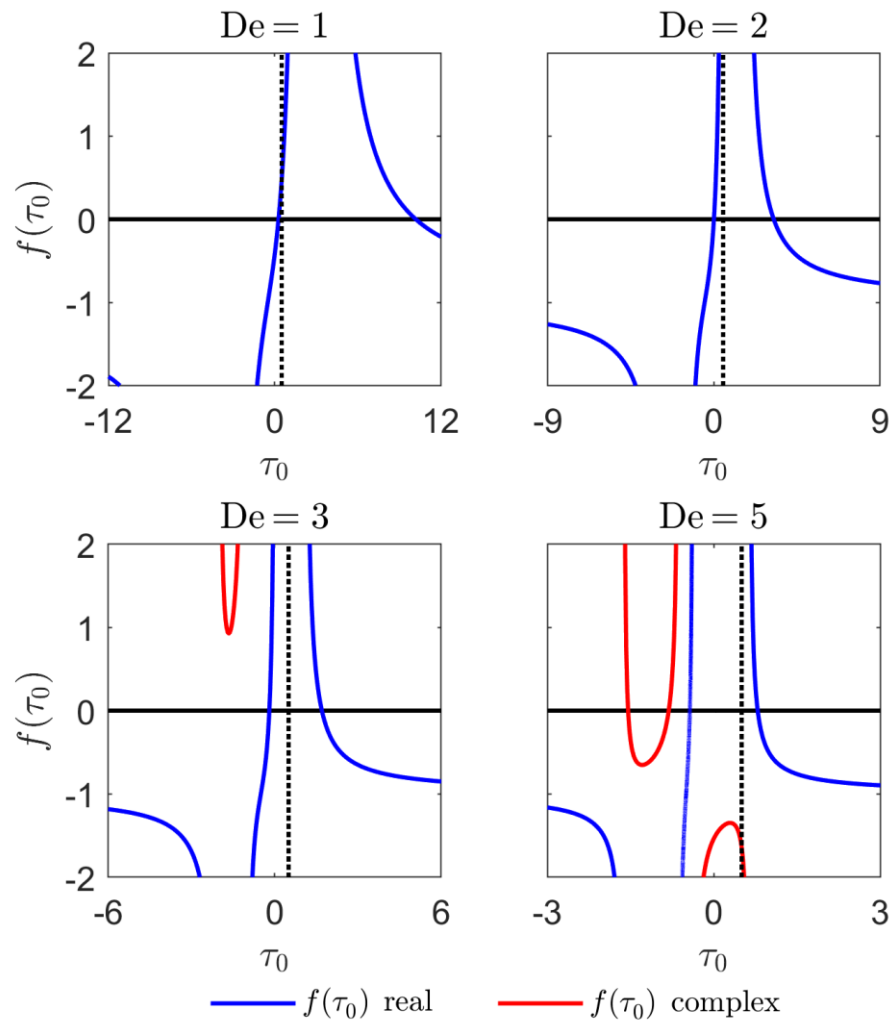


$$u(y) = -\frac{1}{2\alpha G De^2} \left[\frac{2(\alpha - 1)}{1 - \alpha De^2 (\tau_0 + Gy)^2} + (2\alpha - 1) \ln(1 - \alpha De^2 (\tau_0 + Gy)^2) \right] + C.$$

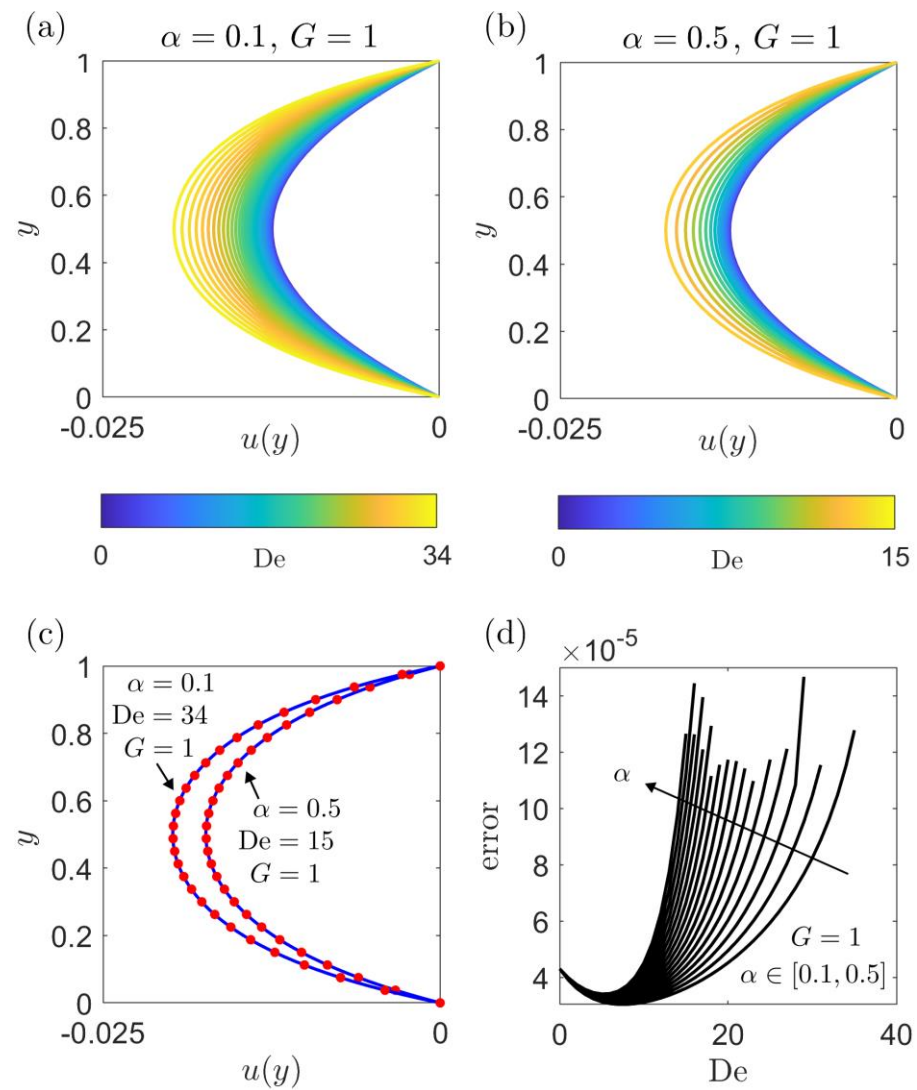
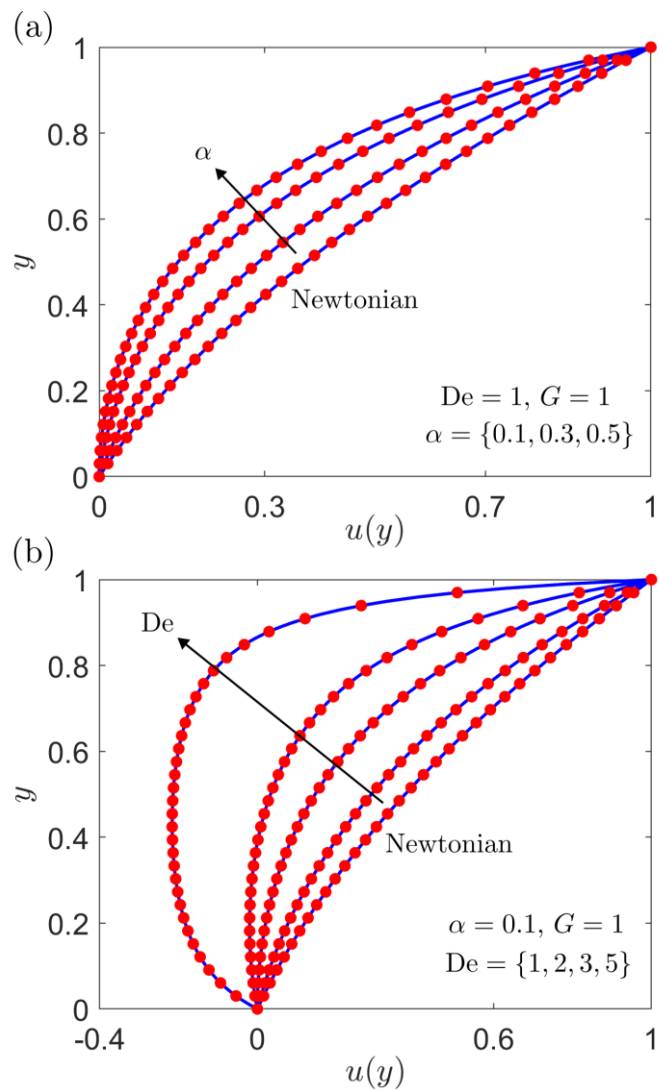
$$\frac{2\alpha - 1}{2\alpha De^2 G} \ln \left(\frac{1 - \alpha De^2 \tau_0^2}{1 - \alpha De^2 (\tau_0 + G)^2} \right) - \frac{(\alpha - 1)(G + 2\tau_0)}{(1 - \alpha De^2 \tau_0^2)(1 - \alpha De^2 (\tau_0 + G)^2)} - 1 = 0$$

A. Raisi. *Rheol. Acta*. 2008.

Extracting the 'physical' flow solution



Comparing the analytical and numerical solutions



Conclusions & future work

Conclusions

- Using a hybrid computational approach, we can efficiently simulate sperm in shear-thinning fluids. A similar technique will allow for the study of bacteria in fluids that exhibit both shear-thinning and viscoelastic properties.
- For sperm cells, shear-thinning rheology tends to hinder propulsion compared to swimming in a Newtonian fluid. This is likely due to flagellar shape changes emerging from fluid-flagellum interactions.
- Even for simple flow problems, obtaining the flow profile of viscoelastic fluids can be difficult. For the Couette-Poiseuille flow of a Giesekus fluid, multiple analytic solutions exist, although our analysis determines that only one solution is physical.

Future work

- To fully implement the hybrid method for the case of modelling bacteria in shear-thinning viscoelastic fluids (and near solid boundaries).
- To explore the effect of non-Newtonian fluids on bacterial locomotion and biofilm development.

Thanks for listening!

Questions?